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A. Wilke, ¹ J. Endres, ² U. Hörmann, ³ J. Niederhausen, ¹ R. Schlesinger, ¹ J. Frisch, ¹ P. Amsalem, ¹ J. Wagner, ³ M. Gruber, ³ A. Opitz, ¹ A. Vollmer, ⁴ W. Brütting, ³ A. Kahn, ² and N. Koch ^{1,4}

¹Institut für Physik, Humboldt-Universität zu Berlin, Brook-Taylor-Str. 15, D-12489 Berlin, Germany ²Department of Electrical Engineering, Princeton University, Princeton, New Jersey 08544, USA

³Institut für Physik, Universität Augsburg, Universitätsstr. 1, D-86135 Berlin, Germany

⁴Helmholtz-Zentrum Berlin für Materialien und Energie GmbH—BESSY II,

Albert-Einstein-Str. 15, D-12489 Berlin, Germany

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We have used ultraviolet and inverse photoemission spectroscopy to determine the transport gaps (E_t) of C_{60} and diindenoperylene (DIP), and the photovoltaic gap (E_{PVG}) of five prototypical donor/acceptor interfaces used in organic photovoltaic cells (OPVCs). The transport gap of C_{60} (2.5 \pm 0.1) eV and DIP (2.55 \pm 0.1) eV at the interface is the same as in pristine films. We find nearly the same energy loss of ca 0.5 eV for all material pairs when comparing the open circuit voltage measured for corresponding OPVCs and E_{PVG} . © 2012 American Institute of Physics. [http://dx.doi.org/10.1063/1.4769360]

The energy level alignment at the donor/acceptor (D/A) heterojunction of an organic photovoltaic cell (OPVC) is decisive for its performance. In particular, the energy offset between the lowest unoccupied molecular orbital level of the acceptor [LUMO(A)] and the highest occupied molecular orbital level of the donor [HOMO(D)] sets an upper limit for the open circuit voltage V_{oc} . This has been expressed as $e \cdot V_{oc} = \text{HOMO(D)} - \text{LUMO(A)} - \Delta$, where e is the elementary charge and Δ a loss term, which has been suggested to be related to the exciton binding energy² or radiative and nonradiative temperature dependent losses. 1,3,5 The HOMO(D)/ LUMO(A) offset is denoted in various ways in the literature, such as charge transfer gap, intermolecular gap, or donor/ acceptor gap and is often estimated by optical spectroscopy, 5,6 or electrical characterization, e.g., cyclic voltammetry,6 reverse saturation current analysis, ^{2,7} or temperature dependent measurements of the open circuit voltage. 4,5 To avoid ambiguity, we use the term *photovoltaic gap* (E_{PVG}) (cf Fig. 1). To minimize energy losses during the photon harvesting process, it is desirable to maximize E_{PVG} within the constraint of keeping the LUMO-LUMO ($\Delta E_{\rm L}$) and HOMO-HOMO level offsets ($\Delta E_{\rm H}$) sufficiently large to drive charge separation across the D/A junction. To quantify Δ for unraveling its physical origin, it is mandatory to have reliable E_{PVG} values for comparison with corresponding V_{oc} values. Unfortunately, simple models for estimating the energy levels at organicorganic interfaces are often invalid (e.g., vacuum level alignment^{8,9}), and more involved models have been brought forward. 10,11 For the time being, experimental determination of interface energetics is indispensable to understand the processes inside an OPVC based on reliable values of E_{PVG} , but only few pertinent studies have been conducted to date. 12-15

The present study focuses on E_{PVG} values at prototypical organic D/A pairs formed between four organic semiconductors [sexithiophene (6T), fullerene (C_{60}), diindenoperylene (DIP, chemical structure shown in Fig. 1), and poly(3-hexylthiophene) (P3HT)] determined using the combination

of ultraviolet and inverse photoemission spectroscopy (UPS and IPES). The experiment for DIP/ C_{60} demonstrates that E_{PVG} can be reliably inferred from measuring the offset between the D/A HOMO levels, once the acceptor's transport gap (E_t) is known and no changes of E_t occur upon interface formation. Further UPS experiments for three other D/A interfaces, i.e., $6T/C_{60}$, 6T/DIP, and P3HT/DIP, and results from previous work for the P3HT/ C_{60} interface interface vield further E_{PVG} values. These are correlated with the V_{oc} of corresponding planar heterojunction (PHJ) solar cells recently investigated, in a reliable value for Δ is obtained.

Substrates consisted of thin films of the commercially available conducting polymer poly(ethylene-dioxythiophene):poly(styrenesulfonate) (PEDT:PSS) (Baytron[®] P AI4083) spin-coated (1500 rpm) on UV/O₃-treated indium tin oxide (ITO) on glass from aqueous dispersion and subsequently annealed at 200 °C for 5 min under ambient conditions. The donor polymer P3HT was spin coated (1500 rpm) on top of PEDT:PSS from a chloroform solution (6 mg/ml) in a N₂-filled glove box and annealed at 180 °C for 30 min. The excess polymer was washed away with the solvent resulting in an insoluble, well defined 3 nm P3HT layer with electronic properties identical to thicker layers. ¹⁹ The deposition of C_{60} , 6T (Sigma Aldrich), and DIP (S. Hirschmann, Univ. Stuttgart, Germany) from resistivity heated crucibles at deposition rates of approximately 0.1 nm/min was done in a preparation chamber (base pressure $<1\times10^{-9}$ mbar) with the substrate at room temperature. The nominal film thickness was monitored with a quartz crystal microbalance.

UPS measurements were carried out at the synchrotron light source BESSY II (Berlin) with 35 eV photons and a hemispherical electron analyzer (Scienta SES 100), and in the Humboldt-Universität laboratory using a He discharge lamp and a Phoibos 100 analyzer. In both systems, the secondary electron cutoff (SECO) spectra were measured with a sample bias of $-10 \, \text{V}$. The error of energy values from UPS is estimated to be smaller than 50 meV. UPS experiments in

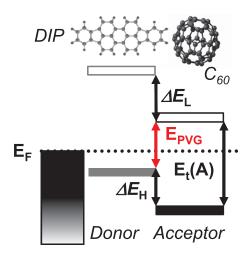


FIG. 1. Scheme of the relevant energy offsets at an organic D/A heterojunction. $\Delta E_{\rm L}$ denotes the LUMO-LUMO offset, $\Delta E_{\rm H}$ the HOMO-HOMO offset. The photovoltaic gap is given by $E_{\rm PVG} = E_{\rm t}({\rm A}) - \Delta E_{\rm H}$. As example, DIP and $C_{\rm 60}$ are shown as donor and acceptor, respectively.

conjunction with IPES measurements were performed at Princeton University using a similar He discharge source and a double-pass cylindrical mirror analyzer. The resolution of all UPS measurements was 150 meV or slightly better.

IPES was performed in the isochromat mode, using a set-up described elsewhere.²⁰ The overall IPES instrumental resolution (450 meV) is estimated from the width of the Fermi step measured on a gold surface. The Gaussian broadening of IPES requires a correction of 50 meV to the onset value of the LUMO and thus to the electron affinity (EA) and the transport gap (E_t) .²¹ The error of energy values from IPES is estimated to be 100 meV.

OPVCs (4 mm² area) were fabricated and characterized at Augsburg University. The open circuit voltage ($V_{\rm OC}$) values were extracted from current-voltage characteristics recorded under white LED (54 mW/cm²) illumination in a N₂-filled glove box at room temperature. The effective intensity is estimated to correspond to 1 sun. To account for the deviation from standard illumination conditions, we assume an error of $V_{\rm OC}$ of ± 15 mV. Temperature variations produce an additional error of ± 5 mV. Thus, the given $V_{\rm OC}$ values are stated within an error of ± 20 mV.

Fig. 2 displays the UPS and IPES spectra of 5 nm C_{60}/ITO (bottom), 10 nm DIP/PEDT:PSS (top), and 1 nm C_{60} on 8 nm DIP/PEDT:PSS (center). All spectra are referenced to the vacuum level (E_{vac}) set to zero. The UPS and IPES spectra are left and right of E_F , respectively. For pristine C_{60} , the HOMO low binding energy (BE) onsets yields an ionization energy (IE) of 6.45 eV, the LUMO high BE onset is at 3.99 eV, and thus an onset-onset gap of 2.46 eV is obtained. Taking into account the experimental Gaussian broadening of IPES results in EA(C₆₀) = 3.94 eV and E_t (C₆₀) = (2.5 ± 0.1) eV, in good agreement with earlier work. 22,23 For the pristine DIP film, the onset procedure results in an IE of 5.35 eV (characteristic of standing DIP molecules in the film¹⁸) and a LUMO high BE onset of 2.85 eV, and hence an onset-onset gap of 2.49 eV. Taking into account the experimental IPES broadening leads to EA(DIP) = $2.80 \,\text{eV}$ and $E_t(\text{DIP}) = (2.55 \,\text{eV})$ \pm 0.1) eV. We attribute the slight deviation from an earlier reported value of 2.7 eV to energetic broadening due to higher

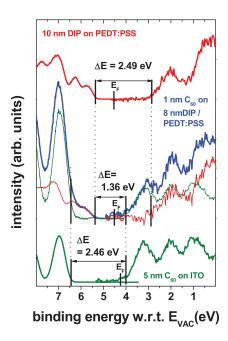


FIG. 2. UPS and IPES spectra of 5 nm C_{60} on ITO (bottom), 1 nm C_{60} on 10 nm DIP/PEDT:PSS (center), and 10 nm DIP on PEDT:PSS (top). Thin lines in die center spectrum are the decomposed contribution of DIP (red) and C_{60} (green). Depicted vertical ticks and dots are guides for the eye.

disorder in our 10 nm DIP film on PEDT:PSS compared to a DIP film grown on Ag(111).²¹

Turning to heterojunctions, we now use the information acquired above for individual materials to assess whether the transport gaps change upon the D/A interface formation with respect to E_t values of the pristine materials. We do so by analyzing the UPS and IPES results from 1 nm C_{60} deposited on 8 nm DIP/PEDT:PSS. Previous thickness-dependent studies of C_{60} on DIP/PEDT:PSS revealed vacuum level alignment and flat bands in both layers, consistent with a weak interaction between the two materials. ¹⁸

We first discuss the UPS results (left of E_F) of the DIP/C₆₀ interface. The C₆₀ spectral contribution (thin green line) at the interface was retrieved by subtracting the unshifted and scaled-down neat DIP UPS valence spectrum (thin red line) measured prior to C₆₀ evaporation from the measured interface spectrum (thick blue line). The subtraction procedure yields "clean" C₆₀ features and no energy shifts upon interface formation, and thus the same C₆₀ and DIP IE values as neat films. The energy offset between the two HOMO levels is 1.1 eV, in agreement with earlier results. ¹⁸

Turning to the IPES results (right of E_F) from the D/A interface, the broadening-corrected onset of the first measured peak gives an EA of 3.94 eV, which is the same value as for a neat C_{60} film. Thus, we attribute the first IPES peak to the C_{60} LUMO. The spectral contribution of DIP at the C_{60} /DIP interface (thin red line) was retrieved by subtracting the appropriately scaled IPES spectrum of the neat C_{60} film (thin green line). The subtraction procedure yields "clean" IPES DIP features, and again no energy shifts upon interface formation. Hence, we find the same C_{60} and DIP EA as in the case of neat C_{60} and DIP films, and an energy offset between the onsets of the DIP and C_{60} LUMO levels equal to 1.15 eV. Finally, these values yield an onset-onset gap of

FIG. 3. (a) SECO and valence region photoemission spectra (synchrotron radiation 35 eV) for increasing C₆₀ coverage on 6T (10 nm)/PEDT:PSS. (b) SECO and valence region photoemission spectra (He II) for increasing DIP coverage on 6T (10 nm)/PEDT:PSS. (c) SECO and (b) valence region photoemission spectra (He I) for increasing DIP coverage on P3HT (3 nm)/PEDT:PSS. Vertical lines and ticks are guides for the eye.

binding energy (eV)

kinetic energy (eV)

1.36 eV and a broadening-corrected E_{PVG} of (1.40 \pm 0.1) eV for the DIP/C₆₀ interface.

The combined IPES and UPS experiments on the DIP/ C_{60} interface show that $E_{\rm t}$ of neither compound changes at the

weakly interacting organic/organic interface with flat bands on either side of the junction. Note that this result, obtained by careful fitting of the spectra, is in contradiction with the result of another UPS/IPES study of the D/A copper-phthalocyanine (CuPc)/ C_{60} interface, where significant changes in polarization energy (and thus in E_t) was suggested for the interface compared to the neat materials. Knowledge of the acceptor E_t at the interface now allows the determination of E_{PVG} of other D/A interfaces by measurement of the energy separation between the D and A HOMO levels at the heterojunction (illustrated in Fig. 1), which we present below.

The UPS results obtained for the C_{60} and DIP grown on 10 nm thick donor 6T layers (on PEDT:PSS) are displayed in Figs. 3(a) and 3(b), respectively. Note that in conjunction with 6T, both C_{60} and DIP take the role of the acceptor. The evolution of the SECO upon deposition of C_{60} on 6T/PEDT:PSS is shown in the left panel of Fig. 3(a). The initial work function of 4.35 eV increases to 4.5 eV upon deposition of 10 nm C_{60} . The small interface dipole of +0.15 eV is most likely due to minute intermolecular charge transfer and/or mutual polarization of molecules at the interface, in analogy to the case of similar heterojunctions. ^{24–26} Note, however, that a clear-cut identification of the origin of the interface dipole at the molecular scale would only be possible with extensive theoretical modeling, which is beyond the scope of this work.

The low BE onset of the 6T HOMO peak is observed $0.35\,\mathrm{eV}$ below $\mathrm{E_F}$ in the valence region shown in the right panel of Fig. 3(a). This yields an IE of $4.7\,\mathrm{eV}$ for the 10 nm thick 6T layer on PEDT:PSS, which corresponds well with previous reports on layers of upright-standing 6T molecules. As long as the 6T levels are still detectable upon increasing $\mathrm{C_{60}}$ coverage, their BE remains unchanged. Applying the subtraction procedure described above to the $1.5\,\mathrm{nm}\,\mathrm{C_{60}}/\mathrm{DIP}$ spectrum yields a "clean" $\mathrm{C_{60}}$ spectrum [thin line in the right panel of Fig. 3(b)] with a low BE HOMO onset at $1.9\,\mathrm{eV}$, which is the same as for larger thicknesses. Consequently, flat band conditions prevail also at the $\mathrm{C_{60}}/\mathrm{6T}$ interface. The energy offset between the low BE onsets of the $\mathrm{C_{60}}$ and 6T HOMO levels is $1.55\,\mathrm{eV}$.

The evolution of the SECO upon deposition of DIP on 6T/PEDT:PSS is shown in the left panel of Fig. 3(b). The initial 6T/PEDT:PSS work function of 4.35 eV is slightly decreased to 4.25 eV upon deposition of 1.2 nm DIP and remains constant for higher DIP coverage. The small interface dipole of -0.1 eV is tentatively attributed to mutual polarization. The thickness-dependent valence spectra of DIP grown on 6T/PEDT:PSS are shown in the right panel of Fig. 3(b). As before, the low BE onset of the 6T HOMO is at 0.35 eV BE. The subtraction procedure reveals the low BE onset of the DIP HOMO emission at 1.1 eV BE [thin line in

TABLE I. HOMO offsets ΔE (HOMO), E_{PVG} , and open circuit voltages (Voc) measured under same illumination conditions at room temperature of investigated D/A interfaces. The transport gap (E_t) of C₆₀ is 2.50 eV and of DIP it is 2.55 eV.

P3HT

D/A pair	$\Delta E(\text{HOMO}) \text{ (eV)}$	E_{PVG} (eV)	Voc (V)	Δ (eV)
DIP/C ₆₀	1.10 ± 0.05	1.40 ± 0.15	0.93 ± 0.02	0.47
6T/C ₆₀	1.55 ± 0.05	0.95 ± 0.15	0.45 ± 0.02	0.50
P3HT/C ₆₀	1.50 ± 0.05	1.00 ± 0.15	0.50 ± 0.02	0.50
6T/DIP	0.75 ± 0.05	1.80 ± 0.15	1.38 ± 0.02	0.42
P3HT/DIP	0.90 ± 0.05	1.65 ± 0.15	1.24 ± 0.02	0.41

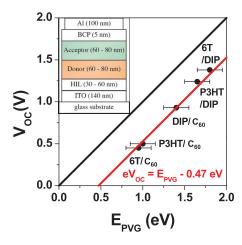


FIG. 4. V_{OC} versus E_{PVG} of planar heterojunction OPVCs with a device structure as shown in the inset.

Fig. 3(b)]. As the layer thickness is increased, no energy shift is detected in the valence region, indicating flat levels in both the DIP and 6T layers. The IE of the DIP layer is 5.35 eV, which corresponds to that of a layer of standing DIP molecules. The energy offset between the low BE emission onsets from the 6T and DIP HOMO levels is 0.75 eV, which is only about half of what was found for 6T/C₆₀.

The interface results for DIP grown on P3HT/PEDT:PSS [Fig. 3(c)] were analyzed in the same way as the previously described interfaces. The experimental results of the D/A HOMO level offsets of DIP and C_{60} on P3HT/PEDT:PSS are given in Table I.

The energy level offsets of the five investigated heterojunctions and the $V_{\rm OC}$ values obtained from the corresponding solar cells are summarized in Table I. In a E_{PVG} vs. V_{OC} plot (Fig. 4), covering an E_{PVG} range from 0.95 eV to 1.80 eV, we find the predicted linear relationship between $V_{\rm OC}$ and E_{PVG} with a slope of unity and an intersection with the energy axis at (0.47 ± 0.05) eV, indicating that for all five OPVCs, the V_{OC} at room temperature is about half a Volt smaller than the HOMO-LUMO energy level offset at the respective D/A interface, i.e., the loss Δ discussed in the Introduction. From the obtained relationship between $V_{\rm OC}$ and $E_{\rm PVG}$, we conclude that a minimum photovoltaic gap of ca 0.5 eV is required in order to obtain a photovoltaic response for planar heterojunction OPVCs at room temperature. Noteworthy, all investigated interfaces exhibit a very similar value of Δ , despite the rather different nature of the two acceptors (DIP vs. C₆₀) and the molecular material 6T vs. the polymer P3HT.

In conclusion, we found the same transport gaps for DIP (2.55 eV) and C_{60} (2.50 eV) at the C_{60} /DIP interface as for the pristine C_{60} and DIP layers. Furthermore, small but finite interface dipoles and flat energy levels were measured for the organic D/A heterojunctions 6T/ C_{60} , 6T/DIP, and P3HT/DIP, but none for DIP/ C_{60} . Comparing the photovoltaic gap of each interface, and in addition also that of P3HT/ C_{60} , with the open circuit voltage measured on corresponding OPVCs, we find that a minimum E_{PVG} of approximately 0.5 eV is necessary for a photovoltaic response of these heterojunctions at room temperature under illumination intensities around 1 sun. The fact that the reliably determined Δ values for vastly dissimilar material pairs are very similar is remar-

kable in the light of the various discussed origins of this energy loss. For instance, it should be important to identify radiative and non-radiative recombination loss mechanisms⁵ at the D/A interfaces and to investigate possible counteracting trends of their ratio.

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¹L. J. A. Koster, V. D. Mihailetchi, R. Ramaker, and P. W. M. Blom, Appl. Phys. Lett. **86**, 123509 (2005).

²B. P. Rand, D. P. Burk, and S. R. Forrest, Phys. Rev. B **75**, 115327 (2007). ³M. Riede, T. Mueller, W. Tress, R. Schueppel, and K. Leo, Nanotechnology **19**, 424001 (2008).

⁴G. Dennler, M. C. Scharber, and C. J. Brabec, Adv. Mater. **21**, 1323 (2009).

⁵M. Gruber, J. Wagner, K. Klein, U. Hörmann, A. Opitz, M. Stutzmann, and W. Brütting, Adv. Energy Mater. 2, 1100 (2012).

⁶D. Veldman, S. C. J. Meskers, and R. A. J. Janssen, Adv. Funct. Mater. **19**, 1939 (2009).

⁷K. Vandewal, K. Tvingstedt, A. Gadisa, O. Inganäs, and J. V. Manca, Phys. Rev. B 81, 125204 (2010).

⁸A. Rajagopal, C. I. Wu, and A. Kahn, J. Appl. Phys. **83**, 2649 (1998).

⁹H. Ishii, K. Sugiyama, W. Ito, and K. Seki, Adv. Mater. **11**, 605 (1999).

¹⁰H. Vazquez, W. Gao, F. Flores, and A. Kahn, Phys. Rev. B **71**, 041306R (2005).

¹¹M. Linares, D. Beljonne, J. Cornil, K. Lancaster, J. L. Brédas, S. Verlaak, A. Mityashin, P. Heremans, A. Fuchs, C. Lennartz, J. Idé, R. Méreau, P. Aurel, L. Ducasse, and F. Castet, J. Phys. Chem. C 114, 3215 (2010).

¹²K. Akaike, A. Opitz, J. Wagner, W. Brütting, K. Kanai, Y. Ouchi, and K. Seki, Org. Electron. 11, 1853 (2010).

¹³Z.-L. Guan, J. B. Kim, Y.-L. Loo, and A. Kahn, J. Appl. Phys. 110, 043719 (2011).

¹⁴E. L. Ratcliff, J. Meyer, K. X. Steirer, N. R. Armstrong, D. Olson, and A. Kahn, Org. Electron. 13, 744 (2012).

¹⁵K. Akaike, K. Kanai, Y. Ouchi, and K. Seki, Adv. Funct. Mater. 20, 715 (2010).

¹⁶F. J. Zhang, A. Vollmer, J. Zhang, Z. Xu, J. P. Rabe, and N. Koch, Org. Electron. 8, 606 (2007).

¹⁷U. Hörmann, J. Wagner, M. Gruber, A. Opitz, and W. Brütting, Phys. Status Solidi (RRL) 5, 241 (2011).

¹⁸J. Wagner, M. Gruber, A. Hinderhofer, A. Wilke, J. Frisch, P. Amsalem, B. Broeker, A. Vollmer, A. Opitz, N. Koch, F. Schreiber, and W. Bruetting, Adv. Funct. Mater. 20, 4295 (2010).

¹⁹J. Frisch, A. Vollmer, J. P. Rabe, and N. Koch, Org. Electron. 12, 916 (2011).

²⁰C. I. Wu, Y. Hirose, H. Sirringhaus, and A. Kahn, Chem. Phys. Lett. 272, 43 (1997).

²¹S. Krause, M. Casu, A. Schöll, and E. Umbach, New J. Phys. **10**, 085001 (2008); S. Krause, doctoral thesis (University Würzburg, 2009), http://opus.bibliothek.uni-wuerzburg.de/volltexte/2009/4047/pdf/Diss_S_Krause.pdf.

²²J. H. Weaver, J. Phys. Chem. Solids **53**, 1433 (1992).

²³R. W. Lof, M. A. van Veenendaal, B. Koopmans, H. T. Jonkman, and G. A. Sawatzky, Phys. Rev. Lett. 68, 3924 (1992).

²⁴S. Verlaak, D. Beljonne, D. Cheyns, C. Rolin, M. Linares, F. Castet, J. Cornil, and P. Heremans, Adv. Funct. Mater. 19, 1 (2009).

²⁵F. Rissner, G. M. Rangger, O. T. Hofmann, A. M. Track, G. Heimel, and E. Zojer, ACS Nano 3, 3513 (2009).

²⁶A. Wilke, P. Amsalem, J. Frisch, B. Bröker, A. Vollmer, and N. Koch, Appl. Phys. Lett. **98**, 123304 (2011).

²⁷S. Duhm, G. Heimel, I. Salzmann, H. Glowatzki, R. L. Johnson, A. Vollmer, J. P. Rabe, and N. Koch, Nature Mater. 7, 326 (2008).