

Rabinowitz Floer homology: a survey

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RABINOWITZ FLOER HOMOLOGY: A SURVEY

PETER ALBERS AND URS FRAUENFELDER

ABSTRACT. Rabinowitz Floer homology is the semi-infinite dimensional Morse homology associated to the Rabinowitz action functional used in the pioneering work of Rabinowitz. Gradient flow lines are solutions of a vortex-like equation. In this survey article we describe the construction of Rabinowitz Floer homology and its applications to symplectic and contact topology, global Hamiltonian perturbations and the study of magnetic fields.

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1. RABINOWITZ FLOER HOMOLOGY

1.1. The Rabinowitz action functional. In his pioneering work [Rab78, Rab79] Rabinowitz used the Rabinowitz action functional to prove existence of periodic orbits on star-shaped hypersurfaces in \mathbb{R}^{2n} . This fundamental work was one of the motivations for Weinstein to state his famous conjecture on the existence of Reeb orbits, [Wei79].

1.1.1. *The unperturbed Rabinowitz action functional and Reeb dynamics.*

Exact symplectic manifolds. Let $(M, \omega = d\lambda)$ be an exact symplectic manifold, for example $(\mathbb{R}^{2n}, \omega_0)$ or a cotangent bundle (T^*B, ω_{std}) each with its canonical symplectic form. We fix an autonomous Hamiltonian, i.e. a smooth time-independent function $F : M \rightarrow \mathbb{R}$. The Hamiltonian vector field X_F of F is defined implicitly by

$$\iota_{X_F}\omega = dF. \quad (1.1)$$

Since F is autonomous the Hamiltonian vector field X_F is tangent to level sets of F and therefore its flow $\phi_F^t : M \rightarrow M$ leaves level sets invariant. This means that the energy F is preserved under the flow ϕ_F^t .

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Let $\mathcal{L} := C^\infty(S^1, M)$, $S^1 = \mathbb{R}/\mathbb{Z}$, be the loop space of M . The Rabinowitz action functional is defined as follows:

$$\begin{aligned} \mathcal{A}^F : \mathcal{L} \times \mathbb{R} &\longrightarrow \mathbb{R} \\ (v, \eta) &\mapsto \mathcal{A}^F(v, \eta) := - \int_{S^1} v^* \lambda - \eta \int_0^1 F(v(t)) dt . \end{aligned} \quad (1.2)$$

The real number η can be thought of as a Lagrange multiplier. Hence critical points $(v, \eta) \in \text{Crit} \mathcal{A}^F$ are critical points of the area functional restricted to the space of loops with F -mean value zero. They are solutions of

$$\begin{cases} \dot{v}(t) = \eta X_F(v(t)), \quad \forall t \in S^1 \\ \int_0^1 F(v(t)) dt = 0 . \end{cases} \quad (1.3)$$

The first equation can be integrated to $v(t) = \phi_F^{\eta t}(v(0))$ and thus, by preservation of energy, the critical point equation is equivalent to

$$\begin{cases} \dot{v}(t) = \eta X_F(v(t)), \quad \forall t \in S^1 \\ v(t) \in F^{-1}(0), \quad \forall t \in S^1 . \end{cases} \quad (1.4)$$

Hence, critical points of \mathcal{A}^F correspond to periodic orbits of X_F with period η and lie on the energy hypersurface $\Sigma := F^{-1}(0)$. Here, the period η is understood in a generalized sense; it is allowed to be negative in which case the periodic orbit is traversed backwards. Moreover, if the period is zero then v is constant and corresponds to a point on the energy hypersurface $F^{-1}(0)$.

If Σ is a regular hypersurface for two functions $F, \tilde{F} : M \rightarrow \mathbb{R}$, $\Sigma = \tilde{F}^{-1}(0) = F^{-1}(0)$, then critical points of $\mathcal{A}^{\tilde{F}}$ agree up to reparametrization with critical points of \mathcal{A}^F . In fact, they are closed characteristics of the canonical line bundle $\ker \omega|_\Sigma \rightarrow \Sigma$ or constant.

It is interesting to compare the critical points of the Rabinowitz action functional \mathcal{A}^F to critical points of the action functional \mathcal{A}_F of classical mechanics

$$\begin{aligned} \mathcal{A}_F : \mathcal{L} &\longrightarrow \mathbb{R} \\ v &\mapsto \mathcal{A}_F(v) := - \int_{S^1} v^* \lambda - \int_0^1 F(v(t)) dt . \end{aligned} \quad (1.5)$$

Critical points $v \in \text{Crit} \mathcal{A}_F$ are 1-periodic solutions of

$$\dot{v}(t) = X_F(v(t)) . \quad (1.6)$$

In this case the period of v is fixed but the energy is arbitrary.

Symplectically aspherical manifolds. A symplectic manifold (M, ω) is called symplectically aspherical if the homomorphism $\bar{I}_\omega : \pi_2(M) \rightarrow \mathbb{R}$ obtained by integrating ω over a smooth representative vanishes identically. This holds for example if $(M, \omega = d\lambda)$ is an exact symplectic manifold. In this situation the Rabinowitz action functional can still be defined on the component $\mathcal{L}_0 \subset \mathcal{L}$ of the loop space of contractible loops. For $v \in \mathcal{L}_0$ there exists a filling disk $\bar{v} : \mathbb{D}^2 \rightarrow M$ where $\mathbb{D}^2 = \{z \in \mathbb{C} \mid |z| \leq 1\}$ and $\bar{v}(e^{2\pi i t}) = v(t)$. In this case we set

$$\begin{aligned} \mathcal{A}^F : \mathcal{L}_0 \times \mathbb{R} &\longrightarrow \mathbb{R} \\ (v, \eta) &\mapsto \mathcal{A}^F(v, \eta) := - \int_{\mathbb{D}^2} \bar{v}^* \omega - \eta \int_0^1 F(v(t)) dt . \end{aligned} \quad (1.7)$$

Due to the symplectic asphericity the definition of \mathcal{A}^F does not depend on the choice of the filling disk. If in addition M is symplectically atoroidal, i.e. $\int_{T^2} f^* \omega = 0$, $\forall f : T^2 \rightarrow M$, the Rabinowitz action functional \mathcal{A}^F can be extended to the whole loop space \mathcal{L} . An interesting class of symplectically atoroidal manifolds are certain twisted cotangent bundles. A twisted cotangent bundle is $(T^*B, \omega_{std} + \tau^* \sigma)$, where $\tau : T^*B \rightarrow B$ is the projection and such that $\sigma \in \Omega^2(B)$ is closed. If the pull-back of σ to the universal cover \tilde{B} has a bounded primitive then $(T^*B, \omega_{std} + \tau^* \sigma)$ is symplectically atoroidal. This fact was used by Merry in [Mer10].

1.1.2. The perturbed Rabinowitz action functional and global Hamiltonian perturbations. Since $F : M \rightarrow \mathbb{R}$ is autonomous the energy hypersurface $\Sigma = F^{-1}(0)$ is preserved under ϕ_F^t . Therefore, Σ is foliated by leaves $L_x := \{\phi_F^t(x) \mid t \in \mathbb{R}\}$, $x \in \Sigma$. It is a challenging problem to compare the system F before and after a global perturbation occurring in the time interval $[0, 1]$. Such a perturbation is described by a function $H : M \times [0, 1] \rightarrow \mathbb{R}$. Moser observed in [Mos78] that it is not possible to destroy all trajectories of the unperturbed system if the perturbation is sufficiently small, that is, there exists $x \in \Sigma$

$$\phi_H^1(x) \in L_x. \quad (1.8)$$

Such a point x is referred to as a leaf-wise intersection point. Equivalently, there exists $(x, \eta) \in \Sigma \times \mathbb{R}$ such that

$$\phi_F^\eta(x) = \phi_H^1(x). \quad (1.9)$$

We point out that the time shift η is uniquely defined by the above equation unless the leaf L_x is closed. If the time shift is negative then the perturbation moves the system back into its own past. Likewise, if the time shift is positive the perturbation moves the system forward into its own future.

Already the existence problem for leaf-wise intersection points is highly non-trivial. The search for leaf-wise intersection points was initiated by Moser in [Mos78] and pursued further in [Ban80, Hof90, EH89, Gin07, Dra08, AF10a, Zil08, AF08, AM09, Gür09, Kan09, AF10b, AF10c, Mer10]. We refer to [AF08] for a brief history.

In [AF10a] Albers – Frauenfelder developed a variational approach to the study of leaf-wise intersection points.

Definition 1.1. A pair $\mathfrak{M} = (F, H)$ of Hamiltonians $F, H : M \times S^1 \rightarrow \mathbb{R}$ is called a Moser pair if it satisfies

$$F(\cdot, t) = 0 \quad \forall t \in [\tfrac{1}{2}, 1] \quad \text{and} \quad H(\cdot, t) = 0 \quad \forall t \in [0, \tfrac{1}{2}], \quad (1.10)$$

and F is of the form $F(x, t) = \rho(t)f(x)$ for some smooth map $\rho : S^1 \rightarrow [0, 1]$ with $\int_0^1 \rho(t)dt = 1$ and $f : M \rightarrow \mathbb{R}$.

If we start with an autonomous Hamiltonian $\hat{F} : M \rightarrow \mathbb{R}$ and an arbitrary $\hat{H} : M \times S^1 \rightarrow \mathbb{R}$ we can find $F, H : M \times S^1 \rightarrow \mathbb{R}$ such that the Hamiltonian flows ϕ_F^t, ϕ_H^t are time reparametrizations of the flows $\phi_{\hat{F}}^t, \phi_{\hat{H}}^t$ and such that (F, H) is a Moser pair.

For simplicity we assume that $(M, \omega = d\lambda)$ is an exact symplectic manifold. For a Moser pair $\mathfrak{M} = (F, H)$ the perturbed Rabinowitz action functional is defined by

$$\begin{aligned} \mathcal{A}^{(F, H)} &\equiv \mathcal{A}^{\mathfrak{M}} : \mathcal{L} \times \mathbb{R} \rightarrow \mathbb{R} \\ (v, \eta) &\mapsto - \int_0^1 v^* \lambda - \int_0^1 H(v, t)dt - \eta \int_0^1 F(v, t)dt. \end{aligned} \quad (1.11)$$

Critical points (v, η) of $\mathcal{A}^{\mathfrak{M}}$ are solutions of

$$\begin{cases} \partial_t v = \eta X_F(v, t) + X_H(v, t), \quad \forall t \in S^1 \\ \int_0^1 F(v, t) dt = 0 \end{cases} \quad (1.12)$$

In [AF10a] we observed that critical points of the perturbed Rabinowitz action functional $\mathcal{A}^{\mathfrak{M}}$ give rise to leaf-wise intersection points.

Proposition 1.2 ([AF10a]). *Let (v, η) be a critical point of $\mathcal{A}^{\mathfrak{M}}$ then $x := v(\frac{1}{2}) \in f^{-1}(0)$ and*

$$\phi_H^1(x) \in L_x \quad (1.13)$$

thus, x is a leaf-wise intersection. Moreover, the map $\text{Crit} \mathcal{A}^{\mathfrak{M}} \rightarrow \{\text{leaf-wise intersections}\}$ is injective unless there exists a leaf-wise intersection point on a closed leaf, see figure 1.

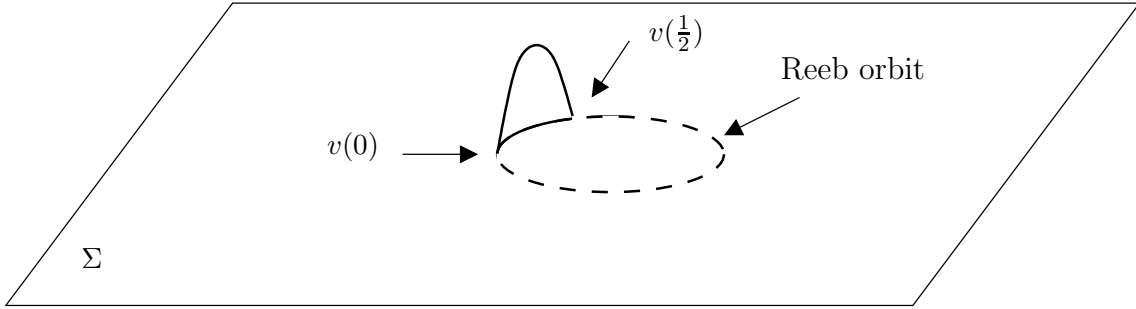


FIGURE 1.

1.1.3. *First properties.* In this paragraph we make the **wrong** assumption that Rabinowitz Floer homology

$$\mathbf{RFH}(\Sigma, M) \equiv \mathbf{HM}(\mathcal{A}^F), \quad (1.14)$$

i.e. the semi-infinite dimensional Morse homology (in the sense of Floer) for the Rabinowitz action functional, can always be constructed and is independent of the defining function F for a regular hypersurface Σ and invariant under Hamiltonian perturbations as described above. Thus, $\mathbf{RFH}(\Sigma, M)$ should have the following properties:

- (1) $\mathbf{RFH}(\Sigma, M) \cong H_*(\Sigma)$ if there are no closed characteristics on Σ . Indeed, then the only critical points of \mathcal{A}^F correspond to constant loops in Σ , see above. Then \mathcal{A}^F is Morse-Bott with $\text{Crit} \mathcal{A}^F \cong \Sigma$ and $\mathcal{A}^F|_{\text{Crit} \mathcal{A}^F} = 0$ since Σ is a regular hypersurface.
- (2) If Σ is displaceable, that is, $\exists H : M \times S^1 \rightarrow \mathbb{R}$ such that $\phi_H^1(\Sigma) \cap \Sigma = \emptyset$, then $\mathbf{RFH}(\Sigma, M) \cong 0$. This follows from invariance under Hamiltonian perturbations

$$\mathbf{RFH}(\Sigma, M) \equiv \mathbf{HM}(\mathcal{A}^F) \cong \mathbf{HM}(\mathcal{A}^{(F,H)}) \cong 0 \quad (1.15)$$

together with the observation that $\text{Crit} \mathcal{A}^{(F,H)} = \emptyset$.

The counterexamples to the Hamiltonian Seifert conjecture, see Ginzburg – Gürel [GG03] and the literature cited therein, are closed hypersurfaces $\Sigma \subset \mathbb{R}^{2n}$ with no closed characteristics. In particular, since any compact subset of \mathbb{R}^{2n} is displaceable we arrive at the contradiction

$$H_*(\Sigma) \cong \mathbf{RFH}(\Sigma, M) \cong 0. \quad (1.16)$$

The reason behind the fact that Rabinowitz Floer homology cannot be defined in full generality is that the moduli spaces of gradient flow lines do not have the necessary compactness properties. However, if certain topological/dynamical assumptions¹ on (Σ, M) are made the desired compactness properties can be established. This is described in the next section.

1.2. Gradient flow lines.

1.2.1. Gradient flow equation. To compute the gradient of the Rabinowitz action functional we need to specify a metric on $\mathcal{L} \times \mathbb{R}$. We take the product of an L^2 -metric on \mathcal{L} and the standard metric on \mathbb{R} . In order to specify the L^2 -metric on \mathcal{L} we choose an S^1 -family $J \equiv \{J_t\}_{t \in S^1}$ of ω -compatible almost complex structures on M . For $(\hat{v}_1, \hat{\eta}_1), (\hat{v}_2, \hat{\eta}_2) \in T_{(v, \eta)}(\mathcal{L} \times \mathbb{R}) = \Gamma(v^*TM) \times \mathbb{R}$ we set

$$\mathbf{m}((\hat{v}_1, \hat{\eta}_1), (\hat{v}_2, \hat{\eta}_2)) \equiv \mathbf{m}_J((\hat{v}_1, \hat{\eta}_1), (\hat{v}_2, \hat{\eta}_2)) := \int_0^1 \omega(\hat{v}_1, J_t(v)\hat{v}_2)dt + \hat{\eta}_1\hat{\eta}_2. \quad (1.17)$$

Then the gradient of \mathcal{A}^F is

$$\nabla^{\mathbf{m}} \mathcal{A}^F(v, \eta) = \begin{pmatrix} J_t(v)(\partial_t v - \eta X_F(v)) \\ - \int_0^1 F(v)dt \end{pmatrix}. \quad (1.18)$$

A gradient flow line is formally a map $w = (v, \eta) \in C^\infty(\mathbb{R}, \mathcal{L} \times \mathbb{R})$ satisfying

$$\partial_s w(s) + \nabla^{\mathbf{m}} \mathcal{A}^F(w(s)) = 0 \quad (1.19)$$

in the sense of Floer, that is $v : \mathbb{R} \times S^1 \rightarrow M$, $\eta : \mathbb{R} \rightarrow \mathbb{R}$ satisfy

$$\begin{cases} \partial_s v + J_t(v)(\partial_t v - \eta X_F(v)) = 0 \\ \partial_s \eta - \int_0^1 F(v)dt = 0. \end{cases} \quad (1.20)$$

If we replace \mathcal{A}^F by the perturbed Rabinowitz action functional $\mathcal{A}^{(F, H)}$ then the corresponding gradient flow equation changes to

$$\begin{cases} \partial_s v + J_t(v)(\partial_t v - \eta X_F(t, v) - X_H(t, v)) = 0 \\ \partial_s \eta - \int_0^1 F(t, v)dt = 0. \end{cases} \quad (1.21)$$

1.2.2. Compactness.

The case of restricted contact type. Let $(W, \omega = d\lambda)$ be a compact, exact symplectic manifold with contact type boundary $\Sigma = \partial W$, that is, the Liouville vector field L , defined by $\iota_L \omega = \lambda$, points outward along Σ . In particular, $(\Sigma, \alpha := \lambda|_\Sigma)$ is a contact manifold.¹ We denote by M the completion of W obtained by attaching the positive half of the symplectization of Σ , that is, $(M = W \cup_\Sigma (\Sigma \times \mathbb{R}_+), \omega = d\lambda)$ where λ is extended over $\Sigma \times \mathbb{R}_+$ by $e^r \alpha$, $r \in \mathbb{R}_+$. We assume that $F : M \rightarrow \mathbb{R}$ is a defining function for Σ such that dF has compact support. The main compactness theorem is as follows:

¹ In general, a closed hypersurface in a symplectic manifold (W, ω) is of *restricted contact type* if there exists a globally defined primitive λ of ω with Liouville vector field L satisfying $L \lrcorner \Sigma$. If λ only exists locally near Σ then Σ is of *contact type*.

Theorem 1.3. *In the above situation let $w_n = (v_n, \eta_n)$ be a sequence of gradient flow lines of $\mathcal{A}^{\mathfrak{M}}$ for which there exists $a < b$ such that*

$$a \leq \mathcal{A}^{\mathfrak{M}}(w_n(s)) \leq b \quad \forall s \in \mathbb{R}. \quad (1.22)$$

Then for every reparametrisation sequence $\sigma_n \in \mathbb{R}$ the sequence $w_n(\cdot + \sigma_n)$ has a subsequence which converges in $C_{\text{loc}}^\infty(\mathbb{R}, \mathcal{L} \times \mathbb{R})$.

The proof follows from standard arguments in Floer theory as soon as we establish

- (1) a uniform $C^0(\mathbb{R})$ bound on η_n ,
- (2) a uniform $C^0(\mathbb{R} \times S^1)$ bound on v_n ,
- (3) a uniform $C^0(\mathbb{R} \times S^1)$ bound on the derivatives of v_n .

Assertion (2) follows from a maximum principle since M is convex at infinity. Assertion (3) follows from standard bubbling-off analysis for holomorphic spheres in Floer theory together with the assumption that (M, ω) is exact. Indeed, a non-constant holomorphic sphere in M has to exist if the derivatives of v_n explode. This contradicts Stokes theorem since ω is exact. Obviously, symplectic asphericity of M would be sufficient.

An interesting feature of gradient flow lines is that they have an infinite amount of flow time but a finite amount of energy available. This leads to the paradoxical conclusion that it is favorable for a gradient flow line to run slowly in order to get far away. The following ‘‘Fundamental Lemma’’ prevents that for slowly running gradient flow lines $w = (v, \eta)$ the Lagrange multiplier η explodes, i.e. slowly running gradient flow lines cannot get too far.

Fundamental Lemma (restricted contact type case). There exists a constant $C > 0$ such that for all $(v, \eta) \in \mathcal{L} \times \mathbb{R}$

$$\|\nabla^{\mathfrak{M}} \mathcal{A}^{\mathfrak{M}}(v, \eta)\| < \frac{1}{C} \implies |\eta| \leq C(|\mathcal{A}^{\mathfrak{M}}(v, \eta)| + 1) \quad (1.23)$$

Remark 1.4. With help of the Fundamental Lemma it can be proved that along gradient flow lines the Lagrange multiplier is bounded in terms of asymptotic action values. This depends heavily on the fact that the action $\mathcal{A}^{\mathfrak{M}}$ is decreasing along gradient flow lines. For details we refer to Cieliebak – Frauenfelder, [CF09a].

Remark 1.5. The assumption of restricted contact type enters the proof of the Fundamental Lemma through the following observation. If we normalize F such that $X_F|_\Sigma$ is the Reeb vector field of α then in the unperturbed case at a critical points $(v, \eta) \in \text{Crit} \mathcal{A}^F$ we have

$$\mathcal{A}^F(v, \eta) = -\eta. \quad (1.24)$$

Since $v(t/\eta)$ is a η -periodic Reeb orbit equation (1.24) can be thought of as a period-action equality for Reeb orbits. Consequently, equation (1.23) can be interpreted as a period-action inequality for almost Reeb orbits.

Remark 1.6. For the unperturbed Rabinowitz action functional the Fundamental Lemma and Theorem 1.3 have been proved in [CF09a] and for the perturbed Rabinowitz action functional in [AF10a].

Remark 1.7. If the defining function F and the almost complex structure J are adapted to the symplectization structure near the hypersurface Σ then the constant C in the Fundamental Lemma can be chosen universally, i.e. independent of F , J , and the contact structure. This has been used in the definition of spectral invariants for Rabinowitz Floer homology in [AF10c].

Remark 1.8. The assumption that the differential of the defining function $F : M \rightarrow \mathbb{R}$ has compact support can be replaced by an appropriate asymptotic growth condition without changing the Rabinowitz Floer homology. This property plays an important role in the proof relating Rabinowitz Floer homology to symplectic homology and loop space topology, see [CFO09, AS09, Mer10].

Remark 1.9. Let (M, ω) be an exact symplectic manifold which is convex at infinity. If $\Sigma \subset M$ is a compact, bounding hypersurface of restricted contact type then the symplectization of Σ embeds into M . It has been examined by Cieliebak – Frauenfelder – Oancea in [CFO09] under which conditions gradient flow lines of the Rabinowitz action functional do not leave the positive part of the symplectization of Σ . The upshot is that M can be replaced by the completion of the region bounded by Σ without changing the Rabinowitz Floer homology.

The case of stable hypersurfaces. Let (M, ω) be a symplectically aspherical symplectic manifold. We assume further that (M, ω) is convex at infinity or geometrically bounded to guarantee C^0 -bounds for gradient flow lines. A closed hypersurface $\Sigma \subset M$ is called stable if there exists $\alpha \in \Omega^1(\Sigma)$ such that

$$\alpha \wedge \omega^{m-1}|_{\Sigma} > 0 \quad \text{and} \quad \ker \omega \subset \ker d\alpha \quad (1.25)$$

where $m = \frac{1}{2} \dim M$. The 1-form α is called a stabilizing 1-form. Σ is called of contact type if $\omega|_{\Sigma} = d\alpha$. Since $\ker \omega|_{\Sigma} \rightarrow \Sigma$ is a rank-1-distribution the Reeb vector field R on Σ is uniquely defined by

$$\iota_R \alpha = 1 \quad \text{and} \quad \iota_R \omega|_{\Sigma} = 0. \quad (1.26)$$

The Rabinowitz action functional \mathcal{A}^F is defined as above. However, the Fundamental Lemma does not carry over to the stable case since already at critical points the period-action equality fails to hold for \mathcal{A}^F in general. We define a modified Rabinowitz action functional by

$$\begin{aligned} \widehat{\mathcal{A}}^F : \mathcal{L} \times \mathbb{R} &\longrightarrow \mathbb{R} \\ (v, \eta) &\mapsto \widehat{\mathcal{A}}^F(v, \eta) := - \int_{S^1} v^* \lambda - \eta \int_0^1 F(v(t)) dt. \end{aligned} \quad (1.27)$$

where $\lambda \in \Omega^1(M)$ satisfies $\lambda|_{\Sigma} = \alpha$. Then the period-action equality holds for $\widehat{\mathcal{A}}^F$. If λ is chosen appropriately then the Fundamental Lemma in the stable case is as follows.

Fundamental Lemma (stable case). There exists a constant $C > 0$ such that for all $(v, \eta) \in \mathcal{L} \times \mathbb{R}$

$$\|\nabla^m \mathcal{A}^F(v, \eta)\| < \frac{1}{C} \quad \implies \quad |\eta| \leq C(|\widehat{\mathcal{A}}^F(v, \eta)| + 1) \quad (1.28)$$

This was proved by Cieliebak – Frauenfelder – Paternain in [CFP09b]. Establishing uniform bounds on the Lagrange multiplier along gradient flow lines in terms of asymptotic action values as described in Remark 1.4 in the stable case involves additional arguments. This is due to the fact that $\widehat{\mathcal{A}}^F$ is not necessarily decreasing along gradient flow lines of \mathcal{A}^F . The crucial observation is that

$$\mathcal{A}_\alpha(v, \eta) := \mathcal{A}^F(v, \eta) - \widehat{\mathcal{A}}^F(v, \eta) = \int v^* \lambda - \int \bar{v}^* \omega : \mathcal{L}_0 \times \mathbb{R} \longrightarrow \mathbb{R} \quad (1.29)$$

is non-increasing along gradient flow lines of \mathcal{A}^F , i.e. \mathcal{A}_α serves as a Lyapunov function for the gradient flow of \mathcal{A}^F . This immediately implies that if w is a gradient flow line of \mathcal{A}^F with

$\lim_{s \rightarrow \pm\infty} w(s) =: w_{\pm} \in \text{Crit} \mathcal{A}^F$ we can estimate

$$\mathcal{A}^F(w_+) - \mathcal{A}_{\alpha}(w_-) \leq \widehat{\mathcal{A}}^F(w(s)) \leq \mathcal{A}^F(w_-) - \mathcal{A}_{\alpha}(w_+) . \quad (1.30)$$

Proving that \mathcal{A}_{α} is non-increasing along gradient flow lines involves the bilinear form

$$\widehat{\mathbf{m}}((\hat{v}_1, \hat{\eta}_1), (\hat{v}_2, \hat{\eta}_2)) := \int_0^1 d\lambda (\hat{v}_1, J_t(v) \hat{v}_2) dt + \hat{\eta}_1 \hat{\eta}_2 . \quad (1.31)$$

for $(\hat{v}_1, \hat{\eta}_1), (\hat{v}_2, \hat{\eta}_2) \in T_{(v, \eta)}(\mathcal{L} \times \mathbb{R}) = \Gamma(v^*TM) \times \mathbb{R}$. We note that in general $\widehat{\mathbf{m}}$ is neither symmetric nor positive definite.

Proposition 1.10. *If F , J , and λ are chosen appropriately, see [CFP09b], the following two properties hold*

$$\begin{aligned} (1) \quad & \mathbf{m} - \widehat{\mathbf{m}} \geq 0 \\ (2) \quad & d\widehat{\mathcal{A}}^F(\cdot) = \widehat{\mathbf{m}}(\nabla^{\mathbf{m}} \mathcal{A}^F, \cdot) . \end{aligned} \quad (1.32)$$

Remark 1.11. Property (2) in Proposition 1.10 can be interpreted as $\nabla^{\widehat{\mathbf{m}}} \widehat{\mathcal{A}}^F = \nabla^{\mathbf{m}} \mathcal{A}^F$ although $\nabla^{\widehat{\mathbf{m}}} \widehat{\mathcal{A}}^F$ is not well-defined since $\widehat{\mathbf{m}}$ might be degenerate.

Remark 1.12. Proposition 1.10 is very sensitive to Hamiltonian perturbations. It is not clear if compactness continues to hold in the perturbed case. However, Kang proved in [Kan09] an analogue of Proposition 1.10 in the case that the hypersurface is of contact type (but not necessarily restricted contact type) and that the Hamiltonian perturbation is localized in the symplectization.

Corollary 1.13. *If F , J , and λ are chosen as in Proposition 1.10 then \mathcal{A}_{α} is a Lyapunov function for the gradient flow of \mathcal{A}^F .*

PROOF. Let w be a gradient flow line of \mathcal{A}^F , i.e. a solution of $\partial_s w(s) + \nabla^{\mathbf{m}} \mathcal{A}^F(w(s)) = 0$. Then we estimate using Proposition 1.10

$$\begin{aligned} \frac{d}{ds} \mathcal{A}_{\alpha}(w(s)) &= d\mathcal{A}^F(w(s))[\partial_s w] - d\widehat{\mathcal{A}}^F(w(s))[\partial_s w] \\ &= d\widehat{\mathcal{A}}^F(w(s))[\nabla^{\mathbf{m}} \mathcal{A}^F] - d\mathcal{A}^F(w(s))[\nabla^{\mathbf{m}} \mathcal{A}^F] \\ &= \widehat{\mathbf{m}}(\nabla^{\mathbf{m}} \mathcal{A}^F, \nabla^{\mathbf{m}} \mathcal{A}^F) - \mathbf{m}(\nabla^{\mathbf{m}} \mathcal{A}^F, \nabla^{\mathbf{m}} \mathcal{A}^F) \\ &\leq 0 . \end{aligned} \quad (1.33)$$

□

1.2.3. Relation to symplectic vortex equations. Symplectic vortex equations have been found independently by Mundet and Salamon and were known in the physics literature as gauged sigma models. They are usually studied for Hamiltonian group actions of compact groups. Symplectic vortex equation with domain the 2-dimensional cylinder can be derived from the gradient flow equation of the moment action functional. We explain in this section how the gradient flow equation for the Rabinowitz action functional can be thought of as the substitute for the (not well-defined) symplectic vortex equations for Hamiltonian \mathbb{R} -actions.

Let (M, ω) be a symplectically aspherical manifold and $F : M \rightarrow \mathbb{R}$ a compactly supported autonomous Hamiltonian function. Then the Hamiltonian flow $\phi_F^t : M \rightarrow M$ can be

thought of as a Hamiltonian \mathbb{R} -action with moment map F . The moment action functional corresponding to this \mathbb{R} -action as introduced by Cieliebak – Gaio – Salamon in [CGS00] is

$$\begin{aligned} \mathbb{A}^F : \mathcal{L}_0 \times C^\infty(S^1, \mathbb{R}) &\longrightarrow \mathbb{R} \\ (v, \eta) &\mapsto \mathbb{A}(v, \eta) := - \int \bar{v}^* \omega - \int \eta(t) F(v(t)) dt. \end{aligned} \quad (1.34)$$

This differs from the Rabinowitz action functional in so far that we have an entire loop $\eta(t)$ of Lagrange multipliers, i.e. there are infinitely many Lagrange multipliers. Critical points of \mathbb{A}^F are solutions of

$$\begin{cases} \partial_t v = \eta(t) X_F(v), \quad \forall t \in S^1 \\ F(v(t)) = 0, \quad \forall t \in S^1. \end{cases} \quad (1.35)$$

The gauge group $\mathcal{H} := C^\infty(S^1, \mathbb{R})$ acts on $\mathcal{L}_0 \times C^\infty(S^1, \mathbb{R})$ by

$$(v, \eta) \mapsto h_*(v, \eta) := (\phi_F^{h(t)}(v(t)), \eta + h'). \quad (1.36)$$

The differential of \mathbb{A}^F is invariant under the action of the gauge group:

$$d\mathbb{A}^F(w)[\hat{w}] = d\mathbb{A}^F(h_* w) \left[dh(w)[\hat{w}] \right] \quad (1.37)$$

for all $w \in \mathcal{L}_0 \times C^\infty(S^1, \mathbb{R})$, $\hat{w} \in T_w(\mathcal{L}_0 \times C^\infty(S^1, \mathbb{R}))$, and $h \in \mathcal{H}$. Since \mathcal{H} is connected we conclude that \mathbb{A}^F is invariant under \mathcal{H} , too. We define the based gauge group

$$\mathcal{H}_0 := \{h \in \mathcal{H} \mid h(0) = 0\}. \quad (1.38)$$

The based gauge group acts freely on $\mathcal{L}_0 \times C^\infty(S^1, \mathbb{R})$. Since for any $\eta \in C^\infty(S^1, \mathbb{R})$ there exists a unique $h_\eta \in \mathcal{H}_0$ such that $\eta + h'_\eta = \int_0^1 \eta(t) dt$ we can identify

$$\begin{aligned} \mathcal{C} : \mathcal{L}_0 \times C^\infty(S^1, \mathbb{R}) / \mathcal{H}_0 &\cong \mathcal{L}_0 \times \mathbb{R} \\ [(v, \eta)] &\mapsto (h_\eta)_*(v, \eta) = \left(\phi_F^{h_\eta(t)}(v(t)), \int_0^1 \eta(t) dt \right). \end{aligned} \quad (1.39)$$

Taking the mean-value of η corresponds to the Coulomb gauge on the trivial \mathbb{R} -bundle over S^1 .

Remark 1.14. It is a remarkable fact that there exists a global slice for the gauge group action. This is related to the fact that \mathbb{R} is abelian. For non-abelian gauge theories in general there exist no global slices. If the moment action functional \mathbb{A}^F is restricted to this global Coulomb slice we obtain the Rabinowitz action functional \mathcal{A}^F . This gives another explanation why a single Lagrange multiplier in the Rabinowitz action functional \mathcal{A}^F eventually leads to a point-wise constraint at critical points (as opposed to an integral constraint).

In order to proceed to the symplectic vortex equations we need to assume that there exists an \mathbb{R} -invariant compatible almost complex structure J on M , i.e.

$$(\phi_F^t)_* J = J \quad \forall t \in \mathbb{R}. \quad (1.40)$$

On compact Lie groups invariant almost complex structure exist by averaging. In the current situation existence is not at all clear but occurs for instance if the flow of F is periodic. Given such an invariant J we define an L^2 inner product on $\mathcal{L}_0 \times C^\infty(S^1, \mathbb{R})$ by

$$\mathfrak{g}((\hat{v}_1, \hat{\eta}_1), (\hat{v}_2, \hat{\eta}_2)) := \int_0^1 \omega(\hat{v}_1, J(v) \hat{v}_2) dt + \int_0^1 \hat{\eta}_1(t) \hat{\eta}_2(t) dt \quad (1.41)$$

for $(\hat{v}_i, \hat{\eta}_i) \in T_{(v, \eta)}(\mathcal{L}_0 \times C^\infty(S^1, \mathbb{R}))$. Then the gradient of \mathbb{A}^F equals

$$\nabla^{\mathfrak{g}} \mathbb{A}^F(v, \eta) = \begin{pmatrix} J(v)(\partial_t v - \eta X_F(v)) \\ F(v) dt \end{pmatrix} \quad (1.42)$$

and the gradient flow equation is

$$\begin{cases} \partial_s v + J(v)(\partial_t v - \eta X_F(v)) = 0 \\ \partial_s \eta - F(v) = 0 \end{cases} \quad (\star)$$

where $v : \mathbb{R} \times S^1 \rightarrow M$ and $\eta : \mathbb{R} \times S^1 \rightarrow \mathbb{R}$ are smooth maps. Since the functional \mathbb{A} and the metric \mathfrak{g} are invariant under $\mathcal{H} = C^\infty(S^1, \mathbb{R})$ so is the gradient flow equation. These are the symplectic vortex equations on the cylinder in temporal gauge. The full symplectic vortex equations are

$$\begin{cases} \partial_s v - \zeta X_F(v) + J(v)(\partial_t v - \eta X_F(v)) = 0 \\ \partial_s \eta - \partial_t \zeta = F(v) \end{cases} \quad (\star\star)$$

where $v : \mathbb{R} \times S^1 \rightarrow M$ and $\eta, \zeta : \mathbb{R} \times S^1 \rightarrow \mathbb{R}$ are smooth maps. This is invariant under the enlarged gauge group $\mathcal{G} := C^\infty(\mathbb{R} \times S^1, \mathbb{R})$ acting by

$$g_*(v, \eta, \zeta) = (\phi_F^g(v), \eta + \partial_t g, \zeta + \partial_s g) . \quad (1.43)$$

Then we have

$$\{(v, \eta) \text{ solves } (\star)\} / \mathcal{H} \cong \{(v, \eta, \zeta) \text{ solves } (\star\star)\} / \mathcal{G} \quad (1.44)$$

since in each \mathcal{G} -orbit $[v, \eta, \zeta]$ there is a representative with $\zeta = 0$ obtained by regauging with $g \in \mathcal{G}$ satisfying $\partial_s g = -\zeta$. The second equation in $(\star\star)$ should be thought of as an equation for the curvature of the connection $\mathbf{A} := \eta dt + \zeta ds$ on the cylinder. Indeed $(\star\star)$ can be written as

$$\begin{cases} \partial_{J, \mathbf{A}} v = 0 \\ * F_{\mathbf{A}}(v) = F(v) \end{cases} \quad (\star\star')$$

where $F_{\mathbf{A}}$ is the curvature of \mathbf{A} and $*$ is the Hodge star operator.

Remark 1.15. We point out that we have two inner products on $\mathcal{L}_0 \times \mathbb{R}$, see diagram (1.45). One is the natural L^2 inner product \mathfrak{m} , see (1.17). The other is $\mathcal{C}^*[\mathfrak{g}]$ where \mathcal{C} is the Coulomb map, see (1.39). We recall that \mathfrak{g} is \mathcal{H}_0 -invariant and we denote by $[\mathfrak{g}]$ the induced metric on the quotient.

$$\begin{array}{ccc} & (\mathcal{L}_0 \times C^\infty(S^1, \mathbb{R}), \mathfrak{g}) & \\ \nearrow \iota & \downarrow \pi & \\ (\mathcal{L}_0 \times \mathbb{R}, \mathfrak{m}) & \xleftarrow{\mathcal{C}} & (\mathcal{L}_0 \times C^\infty(S^1, \mathbb{R}) / \mathcal{H}_0, [\mathfrak{g}]) \end{array} \quad (1.45)$$

We point out that

$$\mathfrak{m} \neq \mathcal{C}^*[\mathfrak{g}] \quad (1.46)$$

which is due to the fact that the infinitesimal gauge action is not \mathfrak{g} -orthogonal to the Coulomb slice. Thus, gradient flow lines of $\nabla^{\mathfrak{m}} \mathcal{A}^F$ are different from gauge orbits of gradient flow lines of $\nabla^{\mathcal{C}^*[\mathfrak{g}]} \mathcal{A}^F$. The latter are the symplectic vortex equations. Therefore, the gradient flow equation for the Rabinowitz action functional is substantially different from the symplectic vortex equations. This also reflected in the respective compactness proofs for gradient flow lines.

1.2.4. Definition of Rabinowitz Floer homology. For a generic Moser pair $\mathfrak{M} = (F, H)$ the perturbed Rabinowitz action functional $\mathcal{A}^{\mathfrak{M}}$ is Morse. Interestingly enough, if Σ is a regular hypersurface, it is enough to perturb H in order to make $\mathcal{A}^{\mathfrak{M}}$ Morse, see [AF10a]. On the other hand, if $H = 0$ then \mathcal{A}^F is never Morse. This is due to the fact that \mathcal{A}^F is invariant under the S^1 -action on the loop space \mathcal{L} . Consequently, S^1 acts on the critical set $\text{Crit}\mathcal{A}^F$. Indeed, if $(v, \eta) \in \text{Crit}\mathcal{A}^F$ is a Reeb orbit then $v(t + \tau, \eta)$, $\tau \in S^1$, is also a critical point. Moreover, $\text{Crit}\mathcal{A}^F$ always contains the constant solutions $(x, 0)$, $x \in \Sigma$. These are the fixed points of the S^1 -action. For generic choice of F the unperturbed Rabinowitz action functional \mathcal{A}^F is Morse-Bott, see [CF09a].

In the following we only give a definition of Rabinowitz Floer homology in the case where $\mathcal{A}^{\mathfrak{M}}$ is Morse. The Morse-Bott case can be defined by choosing an auxiliary Morse function on the critical manifolds and counting gradient flow lines with cascades, see Frauenfelder [Fra04, CF09a]. Moreover, we restrict our attention to $\mathbb{Z}/2$ -coefficients. There's no doubt that Rabinowitz Floer homology can also be defined with \mathbb{Z} -coefficients but so far there is no treatment of coherent orientations for Rabinowitz Floer homology in the literature.

The case of restricted contact type. Let Σ be a hypersurface of restricted contact type and \mathfrak{M} a Moser pair such that $\mathcal{A}^{\mathfrak{M}}$ is Morse. The Rabinowitz Floer chain complex is defined as the Morse complex of $\mathcal{A}^{\mathfrak{M}}$ with help of Novikov sums, see Hofer – Salamon, [HS95], for Floer theory with Novikov rings:

$$\mathbf{RFH}(\mathfrak{M}) := \left\{ \xi = \sum_c \xi_c c \mid \#\{c \in \text{Crit}\mathcal{A}^{\mathfrak{M}} \mid \xi_c \neq 0 \in \mathbb{Z}/2 \text{ and } \mathcal{A}^{\mathfrak{M}}(c) \geq \kappa\} < \infty \forall \kappa \in \mathbb{R} \right\}. \quad (1.47)$$

The boundary operator

$$\partial : \mathbf{RFC}(\mathfrak{M}) \longrightarrow \mathbf{RFC}(\mathfrak{M}) \quad (1.48)$$

is defined by counting gradient flow lines by the standard procedure in Floer homology and satisfies $\partial \circ \partial = 0$ thanks to compactness. Then Rabinowitz Floer homology is

$$\mathbf{RFH}(\mathfrak{M}) := H(\mathbf{RFC}(\mathfrak{M}), \partial). \quad (1.49)$$

The usual theory of continuation homomorphisms in Floer theory implies that Rabinowitz Floer homology is independent of auxiliary data such as the almost complex structure J and the perturbation H . Moreover, if the hypersurface Σ is isotoped through restricted contact type hypersurfaces Rabinowitz Floer homology does not change either.

Remark 1.16. Since Rabinowitz action functional is defined on the full loop space and the differential in the Rabinowitz Floer complex counts topological cylinders we can split Rabinowitz Floer homology into factors labeled by free homotopy classes

$$\mathbf{RFH}(\mathfrak{M}) = \bigoplus_{\gamma \in [S^1, M]} \mathbf{RFH}(\mathfrak{M}, \gamma). \quad (1.50)$$

Definition 1.17. We abbreviate by

$$\mathbf{RFH}(\Sigma, M) := \mathbf{RFH}(\mathfrak{M}, \text{pt}) \quad (1.51)$$

where pt denotes the free homotopy class of a point and

$$\mathbf{RFH}(\Sigma, M) := \mathbf{RFH}(\mathfrak{M}) \quad (1.52)$$

where \mathfrak{M} is a Moser pair such that $\mathcal{A}^{\mathfrak{M}}$ is Morse.

Definition 1.18. Let $(W, \omega = d\lambda)$ be an exact, compact symplectic manifold with contact type boundary $\Sigma := \partial W$. We denote by \widehat{W} the completion of W as described at the beginning of section 1.2.2 and set

$$\begin{aligned} \mathbf{RFH}(\partial W, W) &:= \mathbf{RFH}(\Sigma, \widehat{W}) , \\ \mathbf{RFH}(\partial W, W) &:= \mathbf{RFH}(\Sigma, \widehat{W}) . \end{aligned} \quad (1.53)$$

Theorem 1.19 ([CFO09]). *Let $\Sigma \subset (M, \omega = d\lambda)$ be a restricted contact type closed hypersurface which bounds the compact region W , then Rabinowitz Floer homology does not depend on the exterior $M \setminus W$:*

$$\begin{aligned} \mathbf{RFH}(\Sigma, W) &= \mathbf{RFH}(\Sigma, M) , \\ \mathbf{RFH}(\Sigma, W) &= \mathbf{RFH}(\Sigma, M) . \end{aligned} \quad (1.54)$$

There are two other versions of Rabinowitz Floer homology. For this we fix numbers $a < b$ and define the $\mathbb{Z}/2$ vector space

$$\mathbf{RFC}^{(a,b)}(\mathfrak{M}) := \left\{ \sum_c \xi_c c \mid c \in \text{Crit} \mathfrak{M}, a < \mathcal{A}^{\mathfrak{M}}(c) < b, \xi_c \in \mathbb{Z}/2 \right\} . \quad (1.55)$$

Counting gradient flow lines again defines a differential on $\mathbf{RFC}^{(a,b)}(\mathfrak{M})$. We denote the homology by

$$\mathbf{RFH}^{(a,b)}(\mathfrak{M}) . \quad (1.56)$$

The natural inclusion and projection homomorphisms between $\mathbf{RFC}^{(a,b)}(\mathfrak{M})$ and $\mathbf{RFC}^{(a',b')}(\mathfrak{M})$ turn $\left\{ \mathbf{RFH}^{(a,b)}(\mathfrak{M}) \right\}$ into a bidirected system of vector spaces. Thus, we can take direct and inverse limits:

$$\begin{aligned} \overline{\mathbf{RFH}}(\mathfrak{M}) &:= \varinjlim_{b \rightarrow \infty} \varprojlim_{a \rightarrow -\infty} \mathbf{RFH}^{(a,b)}(\mathfrak{M}) , \\ \underline{\mathbf{RFH}}(\mathfrak{M}) &:= \varprojlim_{a \rightarrow -\infty} \varinjlim_{b \rightarrow \infty} \mathbf{RFH}^{(a,b)}(\mathfrak{M}) . \end{aligned} \quad (1.57)$$

The three vector spaces $\mathbf{RFH}(\mathfrak{M})$, $\overline{\mathbf{RFH}}(\mathfrak{M})$, and $\underline{\mathbf{RFH}}(\mathfrak{M})$ are related by a commutative diagram as follows

$$\begin{array}{ccc} \mathbf{RFH}(\mathfrak{M}) & \xrightarrow{\bar{\rho}} & \overline{\mathbf{RFH}}(\mathfrak{M}) \\ & \searrow \underline{\rho} & \downarrow \kappa \\ & & \underline{\mathbf{RFH}}(\mathfrak{M}) \end{array} \quad (1.58)$$

in which $\bar{\rho}$ is an isomorphism and κ is surjective. We point out that the canonical map $\bar{\rho}$ is not necessarily an isomorphism if \mathbb{Z} -coefficients instead of field coefficients are used. For details and proofs see [CF09b].

Remark 1.20. A hypersurface $\Sigma \subset (M, \omega)$ is of virtual restricted contact type if there exists a cover $\pi : \widetilde{M} \rightarrow M$ and a primitive $\widetilde{\lambda}$ of $\pi^* \omega$ which is bounded on $\pi^{-1}(\Sigma)$ with respect to $\pi^* g$ where g is a Riemannian metric on Σ . The notion of virtual restricted contact type plays an important role in the study of twisted cotangent bundles since energy hypersurfaces above Mañé's critical value are of virtual restricted contact type but in general not of restricted contact type, see [CFP09b] for details. Rabinowitz Floer homology for virtual restricted contact type hypersurfaces can be defined as in the case of restricted contact type on the component of the loop space containing the contractible loops.

The case of stable hypersurfaces. In this section the Rabinowitz action functional is only considered on the space of contractible loops. If restricted contact type is replaced by stability of the hypersurface two difficulties appear. The first is of technical nature namely stability is not an open condition, see [CFP09a]. Therefore, in general it cannot be expected that the unperturbed Rabinowitz action functional can be made Morse-Bott on a given stable hypersurface. It remains true that for a small Hamiltonian perturbation the perturbed Rabinowitz action function is Morse. However, Proposition 1.10 does not allow for Hamiltonian perturbations. Therefore, compactness for families of gradient flow lines might fail as follows: new gradient flow lines from infinity, i.e. with very large Lagrange multipliers, might appear. Nevertheless, for sufficiently small perturbation there is no interaction between gradient flow lines with small and large Lagrange multipliers. To implement this the filtration on the Rabinowitz Floer complex has to be modified to involve both action functionals \mathcal{A}^F and \mathcal{A}_α defined above. Since both action functional are Lyapunov functions along gradient flow lines (see above) we obtain a doubly filtered complex $\text{RFC}^{(a,b),(a',b')}$. Then using a Hamiltonian perturbation of size depending on a, b, a', b' a boundary operator ∂ can still be defined by ignoring gradient flow lines with very large Lagrange multipliers. Moreover, the resulting filtered homology is independent of the small Hamiltonian perturbation. Thus, $\overline{\text{RFH}}$ can be defined as above by taking an inverse-direct limit over a, b, a', b' . We point out that there is no analogue of the definition with Novikov sums for stable hypersurfaces.

The second, more serious difficulty is that the just defined Rabinowitz Floer homology might depend on the stabilizing 1-form $\alpha \in \Omega^1(\Sigma)$ in a very subtle way. Indeed, even though the critical points and the gradient flows do not depend on α Rabinowitz Floer homology $\overline{\text{RFH}}$ might depend on α through the double action filtration by \mathcal{A}^F and \mathcal{A}_α in the inverse-direct limit. To guarantee independence of Rabinowitz Floer homology we impose the additional assumption of tameness.

Definition 1.21. A stable hypersurface $\Sigma \in (M, \omega)$ is tame if there exists a stabilizing 1-form $\alpha \in \Omega^1(\Sigma)$ and a constant $C > 0$ such that

$$\left| \int v^* \alpha \right| \leq C \left| \int \bar{v}^* \omega \right| \quad (1.59)$$

for all Reeb orbits v of α which are contractible in M . Here \bar{v} again denotes a filling disk and the right hand side is independent of \bar{v} since (M, ω) is symplectically aspherical.

There are examples of stable hypersurfaces which are not tame. In fact, even contact type hypersurface need not be tame. For more details we refer to [CFP09b]. More examples of non-tame hypersurfaces are provided in Macarini – Paternain [MP09] and Cieliebak – Volkov [CV10]. An isotopy of hypersurfaces is called a stable isotopy if the stabilizing 1-form depends smoothly on the isotopy parameter. It is called a tame stable isotopy if in addition all hypersurfaces are tame with a taming constant independent of the isotopy parameter.

Theorem 1.22 ([CFP09b]). *If Σ is tame then Rabinowitz Floer homology $\overline{\text{RFH}}$ is invariant under tame stable isotopies. Since the space of stabilizing 1-forms is a cone, hence connected, $\overline{\text{RFH}}$ is independent of α .*

Definition 1.23. We abbreviate

$$\text{RFH}(\Sigma, M) := \overline{\text{RFH}}. \quad (1.60)$$

This is justified since in the case of restricted contact type the natural map $\bar{\rho}$ in the commuting diagram 1.58 is an isomorphism.

Grading. If the homomorphism $I_{c_1(M)} : \pi_2(M) \rightarrow \mathbb{Z}$ obtained by integrating the first Chern class $c_1(M)$ over a smooth representative vanishes identically then Rabinowitz Floer homology RFH_* carries a \mathbb{Z} -grading in terms of the Conley-Zehnder index, see [CF09a]. In general only a relative \mathbb{Z}/N can be defined where N is the minimal Chern number. There are different normalization conventions. In [CF09a] the grading takes values in the set $\frac{1}{2} + \mathbb{Z}$. Underlying this convention is the philosophy that the index in a semi-infinite homology is $\frac{1}{2}$ times some kind of signature index of the Hessian at critical points. We refer to the article [RS96] by Salamon – Robbin for research in this direction. The half signature index takes values in \mathbb{Z} on even dimensional manifolds and in $\frac{1}{2} + \mathbb{Z}$ on odd dimensional manifolds. Since the loop space \mathcal{L} of a symplectic manifold is itself symplectic \mathcal{L} is an “even” infinite-dimensional manifold and consequently $\mathcal{L} \times \mathbb{R}$ is an “odd” infinite-dimensional manifold. Therefore the convention used in [CF09a] is consistent with this interpretation. Moreover, with this convention

$$\text{RFH}^{-*} \cong \text{RFH}_* \quad (1.61)$$

holds. This isomorphism is induced by the involution

$$\begin{aligned} I : \mathcal{L} \times \mathbb{R} &\longrightarrow \mathcal{L} \times \mathbb{R} \\ (v, \eta) &\mapsto (v^-, -\eta) \end{aligned} \quad (1.62)$$

where $v^-(t) := v(1 - t)$ under which the unperturbed Rabinowitz action functional is anti-invariant

$$\mathcal{A}^F \circ I = -\mathcal{A}^F. \quad (1.63)$$

Another useful convention is to replace the above convention by adding $\frac{1}{2}$. This convention fits better with the computational results comparing Rabinowitz Floer homology with symplectic homology and loop space topology as carried out in [CFO09, AS09, Mer10].

1.3. Computations.

1.3.1. *The displaceable case.* Let (M, ω) be symplectically aspherical, and convex at infinity or geometrically bounded and $\Sigma \subset (M, \omega)$ of (virtual) restricted contact type or stable tame.

Theorem 1.24.

- (1) *If the hypersurface Σ is displaceable then Rabinowitz Floer homology $\text{RFH}_*(\Sigma, M)$ vanishes.*
- (2) *If in addition Σ is of restricted contact type then even the full Rabinowitz Floer homology $\mathbf{RFH}_*(\Sigma, M)$ vanishes.*

This is proved for restricted contact type in [CF09a] and for the remaining cases in [CFP09b].

Theorem 1.25. *If the hypersurface Σ carries no Reeb orbits which are contractible in M then*

$$\text{RFH}_*(\Sigma, M) \cong H_*(\Sigma). \quad (1.64)$$

This follows immediately from the definition of RFH_* in the Morse-Bott case, see above. The previous two Theorems imply the following Corollary.

Corollary 1.26. *If the hypersurface Σ is displaceable then there exists a Reeb orbit which is contractible in M .*

Remark 1.27. If the hypersurface Σ is stable then Corollary 1.26 has been proved by Schlenk in [Sch06] using different methods. Schlenk does not need the tameness assumption. Using a local version of Rabinowitz Floer homology Corollary 1.26 can also be proved without the tameness assumption in the framework of Rabinowitz Floer homology, see [CFP09b].

1.3.2. Relations to symplectic homology and loop space topology. For the next theorem we assume that we are in the same setup as at the beginning of section 1.2.2, *The case of restricted contact type*. Namely, $(W, \omega = d\lambda)$ is a compact exact symplectic manifold with boundary $\partial W = \Sigma$ of restricted contact type and (M, ω) is the completion of W . In this situation symplectic homology $\mathrm{SH}_*(\Sigma, M)$ resp. cohomology $\mathrm{SH}^*(\Sigma, M)$ of M can be defined, see Cieliebak – Floer – Hofer [CFH95] and Viterbo [Vit99].

Theorem 1.28 ([CFO09]). *There exists a long exact sequence*

$$\cdots \rightarrow \mathrm{SH}^{-*}(\Sigma, M) \rightarrow \mathrm{SH}_*(\Sigma, M) \rightarrow \mathbf{RFH}_*(\Sigma, M) \rightarrow \mathrm{SH}^{-*+1}(\Sigma, M) \rightarrow \cdots \quad (1.65)$$

From now on $M = T^*B$ is the cotangent bundle with its standard symplectic structure $\omega_{std} = d\lambda_{std}$ over a closed manifold B and $\Sigma = S_g^*B$ is the unit cotangent bundle with respect to some Riemannian metric g on B . Using that the symplectic (co-)homology of cotangent bundles has been computed before by Viterbo, Salamon – Weber, and Abbondandolo – Schwarz [Vit99, SW06, AS06] the following Theorem has been proved via the long exact sequence in [CFO09]. An independent and direct proof has been given in [AS09].

Theorem 1.29 ([CFO09, AS09]).

$$\mathbf{RFH}_*(S_g^*B, T^*B) = \begin{cases} \mathrm{H}^{-*+1}(\mathcal{L}_B) & \text{if } * < 0 \\ \mathrm{H}_*(\mathcal{L}_B) & \text{if } * > 1 \end{cases} \quad (1.66)$$

where $\mathcal{L}_B := C^\infty(S^1, B)$ is the free loop space of B .

Remark 1.30. In the remaining degrees $* = 0, 1$ the answer is known and depends on the Euler class, see [CFO09, AS09].

Remark 1.31. In [Mer10] Merry extends Theorem 1.29 to energy hypersurfaces above the Mañé critical value in twisted cotangent bundles.

2. APPLICATIONS AND RESULTS

2.1. Symplectic and Contact topology. As we pointed out in Theorem 1.19 Rabinowitz Floer homology does not depend on the exterior. It is an open question to what extent Rabinowitz Floer homology $\mathrm{RFH}(\Sigma, M)$ depends on the filling M . A partial independence result is the following Theorem.

Theorem 2.1. *Let $\dim B \geq 4$ and $\pi_1(B) = \{1\}$ and let $(W, \omega = d\lambda)$ be a compact exact symplectic manifold with $\partial W \cong S_g^*B$. If $(\partial W, \lambda|_{\partial W})$ is contactomorphic to $(S_g^*B, \lambda_{std}|_{S_g^*B})$ then*

$$\mathbf{RFH}_*(\partial W, W) \cong \mathbf{RFH}_*(S_g^*B, T^*B). \quad (2.1)$$

This is a special case of a result proved in [CFO09]. The crucial ingredient in the above Theorem is that there exist no rigid finite energy planes in the filling W . If B is a sphere then the above theorem can in fact be checked by a direct computation, see [CF09a].

Corollary 2.2. *Under the assumptions of Theorem 2.1 S_g^*B does not admit an exact contact embedding into \mathbb{R}^{2n} or, more generally, into a subcritical Stein manifold.*

PROOF. We assume by contradiction that there exists an exact contact embedding of S_g^*B into a subcritical Stein manifold $(M, \omega = d\lambda)$. We denote by Σ the image of S_g^*B in M . Since $H_{2n-1}(M) = 0$, $2n = \dim M$, the hypersurface Σ bounds a compact region W in M . Because any compact subset of a subcritical Stein manifold is displaceable, see Biran – Cieliebak [BC02], $\mathbf{RFH}_*(\Sigma, M) \cong 0$. On the other hand by Theorem 1.19 and 2.1 we know $0 \cong \mathbf{RFH}_*(\Sigma, M) \cong \mathbf{RFH}_*(\Sigma, W) \cong \mathbf{RFH}_*(S_g^*B, T^*B)$. This contradicts Theorem 1.29. \square

2.2. Global perturbations of Hamiltonian systems. We recall from section 1.1.2 that critical points of $\mathcal{A}^{\mathfrak{M}}$, $\mathfrak{M} = (F, H)$, give rise to leaf-wise intersections.

Theorem 2.3. *Let $\Sigma \subset (M, \omega = d\lambda)$ be a closed, bounding, restricted contact type hypersurface and M convex at infinity or geometrically bounded. We denote by $\wp > 0$ the smallest period of a Reeb orbit on $(\Sigma, \lambda|_{\Sigma})$ which is contractible in M . Let $\psi \in \text{Ham}_c(M, \omega)$, that is, ψ is a Hamiltonian diffeomorphism generated by a Hamiltonian with compact support. If the Hofer norm $\|\psi\| < \wp$ then there exists a leaf-wise intersection.*

This was proved in [AF10a]. An alternative proof is given by Gürel in [Gür09]. Kang [Kan09] found an extension of Theorem 2.3 if Σ is only of contact type. In [AM10] a cup-length estimate for leaf-wise intersections is proved under the assumptions of Theorem 2.3.

Remark 2.4. The existence of displaceable hypersurfaces Σ shows that a smallness assumption $\|\psi\| < \wp$ is necessary.

Remark 2.5. If $\psi = \text{id}$ in Theorem 2.3 then each point in Σ is a leaf-wise intersection point via constant critical points. The leaf-wise intersection point found in Theorem 2.3 arises as a perturbation of a constant critical point by a stretching-of-the-neck argument for gradient flow lines. If this stretching-of-the-neck argument is replaced by local Floer homology around the constant critical points then for a generic Hamiltonian diffeomorphism ψ in Theorem 2.3 there exists $\sum_{i=0}^{\dim \Sigma} b_i(\Sigma)$ different leaf-wise intersection points, see [AF10a, Kan09].

We recall that $\mathcal{L}_B = C^\infty(S^1, B)$ denotes the free loop space of B .

Theorem 2.6. *Let $\dim H_*(\mathcal{L}_B) = \infty$. If $\dim B \geq 2$ and $\Sigma \subset T^*B$ is a generic fiber-wise star-shaped hypersurface then for a generic $\psi \in \text{Ham}_c(T^*B)$ there exist infinitely many leaf-wise intersection points.*

We point out that there is no assumption on the Hofer norm $\|\psi\|$ of ψ in Theorem 2.6. This was proved in [AF08] along the following lines. Since Σ is fiber-wise star-shaped it is of restricted contact type and isotopic to S_g^*B through restricted contact type hypersurfaces. Thus, by the assumption $\dim H_*(\mathcal{L}_B) = \infty$, invariance of Rabinowitz Floer homology, and Theorem 1.29 we conclude that the Rabinowitz Floer homology $\mathbf{RFH}(\Sigma, T^*B)$ is infinite dimensional. In particular, since for a generic perturbation the corresponding Rabinowitz action functional $\mathcal{A}^{\mathfrak{M}}$ is Morse, $\mathcal{A}^{\mathfrak{M}}$ has infinitely many critical points. Then a transversality result using $\dim B \geq 2$ yields that generically critical points of $\mathcal{A}^{\mathfrak{M}}$ won't lie on a closed leaf. Thus, the assertion of Theorem 2.6 follows from Proposition 1.2.

Remark 2.7. Theorems 2.3 and 2.6 have been proved for energy hypersurfaces above Mañé's critical value in twisted cotangent bundles in [Mer10].

Remark 2.8. As the sketch of the proof of Theorem 2.6 shows to obtain generically infinitely many leaf-wise intersection points it is enough to prove that Rabinowitz Floer homology is infinite dimensional. A large class of examples with this property has been constructed by Albers – McLean in [AM09].

Using spectral invariants in Rabinowitz Floer homology Theorem 2.6 can be improved as follows, see [AF10c].

Theorem 2.9. *Let $\dim H_*(\mathcal{L}_B) = \infty$. If $\dim B \geq 2$ and $\Sigma \subset T^*B$ is a fiber-wise star-shaped hypersurface then for $\psi \in \text{Ham}_c(T^*B)$ there exist infinitely many leaf-wise intersection points or a leaf-wise intersection point on a closed leaf.*

This Theorem is equivalent to the assertion that the Rabinowitz action functional always, even in degenerate situations, has infinitely many critical points. This does not follow from Theorem 2.6 since for degenerate functions the Morse inequalities might fail.

The new ingredient in Theorem 2.9 are spectral invariants. The idea behind spectral invariants is the following. To each homology class a critical value is associated. If the action functional is Morse this is done by a mini-max procedure. It can then be shown that this assignment is locally Lipschitz continuous in changes of the Moser pair. Thus, it extends to all Moser pairs. The crucial property of the extension is that it still assigns critical values to homology classes, even in degenerate situations.

To prove Theorem 2.9 it is shown that the set of spectral values is unbounded and therefore gives rise to infinitely many critical points which are distinguished by their critical values.

It is an interesting open question whether a Gromoll-Meyer type theorem, [GM69, Mat80], holds for leaf-wise intersection points, that is, if there exist infinitely many leaf-wise intersection points even in the case when there is a leaf-wise intersection point on a closed leaf. This problem is intimately related to the existence problem of (geometrically distinct) geodesics. Katok's example, [Kat73] suggests that, as in the Gromoll-Meyer Theorem, a growth condition for the homology of the loop space is necessary. Interesting research in the direction can be found in [GG07].

Since $S^{2n-1} \subset \mathbb{R}^{2n}$ is displaceable for general Hamiltonian perturbations there need not exist leaf-wise intersection points. However, if the class of Hamiltonian perturbations is restricted to preserve symmetries this picture might change dramatically. Such a phenomenon was discovered by Ekeland – Hofer in [EH89]. They prove that if Σ is a centrally symmetric, restricted contact type hypersurface in \mathbb{R}^{2n} and ϕ^t is an isotopy of centrally symmetric Hamiltonian perturbations there always exists a leaf-wise intersection point for ϕ^1 . If restricted contact type is replaced by star-shaped it is proved in [AF10b] that under the same symmetry assumptions there exist infinitely many leaf-wise intersection points or a leaf-wise intersection point on a closed leaf. The proof uses a computation of $\mathbb{Z}/2$ -invariant Rabinowitz Floer homology.

2.3. Mañé's critical value. Let (B, g) be a closed Riemannian manifold and $\sigma \in \Omega^2(B)$ be a closed 2-form which vanishes on $\pi_2(B)$. Thus, on the universal cover $\pi : \tilde{B} \rightarrow B$ the 2-form $\pi^*\sigma$ is exact. The twisted cotangent bundle is $(T^*B, \omega := \omega_{std} + \tau^*\sigma)$ where $\tau : T^*B \rightarrow B$ is the projection. We fix a potential $U : B \rightarrow \mathbb{R}$ and define $F : T^*B \rightarrow \mathbb{R}$ by $F(q, p) := \frac{1}{2}\|p\|_g^2 + U(q)$. Then the Mañé critical value is

$$c = c(g, \sigma, U) := \inf_{\theta} \sup_{q \in \tilde{B}} \tilde{F}(q, \theta_q) \quad (2.2)$$

where $\tilde{F} = F \circ \pi$ and $\theta \in \Omega^1(\tilde{B})$ satisfies $d\theta = \pi^*\sigma$.

In physical terms the Hamiltonian dynamics of X_F taken with respect to the twisted symplectic form describes the motion of a particle on B subject to a conservative force $-\nabla U$

and a magnetic field σ . For energies above the Mañé critical value the magnetic field is interpreted as small compared to the energy.

The hypersurfaces $\Sigma_k := F^{-1}(k)$ above the Mañé critical value, i.e. $k > c$, are of virtual restricted contact type, see Remark 1.20. If there are no topological obstructions Σ_k is displaceable for sufficiently small k if $\sigma \neq 0$. This follows from the fact that the zero-section B in T^*B is displaceable, see Polterovich [Pol95].

Theorem 2.10 ([Mer10]). *The hypersurface Σ_k for $k > c$ is not displaceable.*

This follows from the fact that Rabinowitz Floer homology above the Mañé critical value does not vanish, see Remark 1.31 .

At the Mañé critical value the hypersurface Σ_c is in general not stable or of virtual restricted contact type. It might even be singular, e.g. if the magnetic field is zero. Moreover, there are examples where Σ_c carries no closed characteristics, e.g. the horocycle flow. The article [CFP09b] proposes the following *paradigms*:

- ($k > c$) Above the Mañé critical value, Σ_k is virtually contact. It may or may not be stable. Its Rabinowitz Floer homology $\text{RFH}(\Sigma_k)$ is defined and nonzero, so Σ_k is non-displaceable (now proved by Merry [Mer10]). The dynamics on Σ_k is like that of a geodesic flow; in particular, it has a periodic orbit in every nontrivial free homotopy class.
- ($k = c$) At the Mañé critical value, Σ_k is non-displaceable and can be expected to be non-stable (hence non-contact).
- ($k < c$) Below the Mañé critical value, Σ_k may or may not be of contact type. It is stable and displaceable (provided that $\chi(M) = 0$), so its Rabinowitz Floer homology $\text{RFH}(\Sigma_k)$ is defined and vanishes. In particular, Σ_k has a contractible periodic orbit.

It is unknown whether Σ_k is always stable below the Mañé critical value. In [CFP09b] many examples can be found supporting these paradigms.

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