



# TLS Kondo effect in structurally disordered ThAsSe

T. Cichorek, Z. Henkie, J. Custers, Philipp Gegenwart, F. Steglich

### Angaben zur Veröffentlichung / Publication details:

Cichorek, T., Z. Henkie, J. Custers, Philipp Gegenwart, and F. Steglich. 2004. "TLS Kondo effect in structurally disordered ThAsSe." *Journal of Magnetism and Magnetic Materials* 272-276: 66–67. https://doi.org/10.1016/j.jmmm.2003.12.1109.





## TLS Kondo effect in structurally disordered ThAsSe

T. Cichorek<sup>a,b,\*</sup>, Z. Henkie<sup>b</sup>, J. Custers<sup>a</sup>, P. Gegenwart<sup>a</sup>, F. Steglich<sup>a</sup>

<sup>a</sup> Max Planck Institute for Chemical Physics of Solids, Dresden D-01187, Germany
<sup>b</sup> Institute of Low Temperature and Structure Research, Polish Academy of Sciences, 50-950 Wrocław, Poland

#### Abstract

Low-temperature electrical-resistivity,  $\rho(T)$ , measurements on single crystals of the structurally disordered ThAsSe reveal an anomalous scattering mechanism, which is apparently derived from two-level systems. For the ThAsSe specimen displaying a resistivity saturation at millikelvin temperatures, a crossover from a logarithmic to a non-Fermiliquid behavior  $\Delta\rho \propto 1 - aT^{1/2}$  is observed upon cooling below  $\simeq 4$  K. A comparison of experimental results with the theoretical ones yields a characteristic energy scale of the order of a few K for ThAsSe.

Non-magnetic interaction of a degenerate Fermi gas with structural two-level systems (TLS) may cause a logarithmic correction to  $\rho(T)$ , i.e., a behavior analogous to the spin Kondo effect [1]. Even more remarkable is a realization of a two-channel Kondo effect, as predicted when the Kondo temperature,  $T_{\rm K}$ , exceeds the energy splitting between the localized states,  $\Delta_{\rm TLS}$ . A temperature variation of the additional resistivity due to the TLS–electron interaction,  $\rho_{\rm TLS}$ , is shown in the inset of Fig. 1 in a schematic way: Upon cooling below  $T_{\rm K}$ , the  $-\log T$  dependence transforms into the non-Fermi liquid (nFL) behavior  $\rho_{\rm TLS} \propto 1 - a T^{1/2}$ . At the lowest temperatures,  $\rho_{\rm TLS}$  saturates and the Fermi-liquid state eventually develops [1].

Recently, we have reported an anomalous scattering mechanism in the structurally disordered ThAsSe [2]. At  $T \lesssim 20$  K, its  $\rho(T)$  displays a  $-\log T$  correction that holds over one decade in temperature and is unaffected by neither magnetic fields up to 13.5 T (cf. the inset of Fig. 2) nor hydrostatic pressures as high as 1.88 GPa [3]. Additionally, for various single crystals a saturation of

 $\rho(T)$  between 0.2 and 2 K was observed [2]. Most probably, the low-lying excitations of the electron gas in ThAsSe are due to the TLS–electron interaction. Dynamic disorder in this diamagnet was reflected by a quasilinear-in-T term of non-electronic origin in the specific heat at  $T\lesssim 1.7$  K [2].

An example of a complex  $\rho(T)$  behavior in ThAsSe, presented as the  $\Delta \rho(T)/\rho(0.05~{\rm K})$  data on a semi-logarithmic temperature scale, is given in Fig. 1. Following observations are made below 20 K: (i) At higher temperatures,  $\rho(T)$  logarithmically increases with lowering temperature down to  $T^*$ . A rough approximation for  $T^*$  yields a value somewhat larger than 2 K. (ii) Upon further cooling,  $\rho(T)$  significantly deviates from the  $-\log T$  dependence. (iii) An increase of  $\rho(T)$  holds down to  $T_S \simeq 0.2~{\rm K}$ , at which the saturation sets in.

In order to determine a temperature variation of the electrical resistivity at  $T_{\rm S} < T < T^*$ , we re-plotted our data on a  $T^{1/2}$  scale. The results are depicted in Fig. 2, where the value of  $\rho_{ab}(300~{\rm K})=220~\mu\Omega$  cm was taken from Ref. [4]. As indicated by the solid line, a  $T^{1/2}$  dependence holds between around  $\sim 0.2$  and 6 K. The latter value is larger than  $T^*=2~{\rm K}$ , as estimated from Fig. 1. We emphasize that an unambiguous approximation of  $T^*$  cannot be expected in real, macroscopic samples where there are many TLS centers [1].

<sup>\*</sup>Corresponding author. Max Planck Institute for Chemical Physics of Solids, Noethnizer Str. 40, Dresden D-01187, Germany. Tel.: +49-351-4646-3133; fax: +49-351-4646-3119. E-mail address: cichorek@cpfs.mpg.de (T. Cichorek).

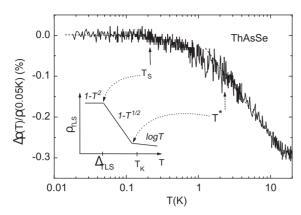


Fig. 1. The *ab*-plane resistivity normalized to the corresponding value at 0.05 K,  $\Delta \rho(T)/\rho(0.05 \text{ K})$ , for the ThAsSe single crystal. The solid arrow at  $T_{\rm S}$  marks the temperature of saturation. The dashed arrow indicates the temperature,  $T^*$ , below which  $\rho(T)$  deviates from a  $-\log T$  dependence. Dashed lines are a guide to the eye only. Inset: Expected temperature dependence of  $\rho_{\rm TLS}$  shown in a schematic way for the case of  $\Delta_{\rm TLS} \ll T_{\rm K}$ .

As mentioned above, the anomalous low-T properties of ThAsSe are apparently caused by the TLS-electron interaction. Furthermore, the ThAsSe sample with  $T_{\rm S} \simeq 0.2 \, \rm K$  displays a far-reaching similarity of the experimental  $\rho(T)$  results to the theoretical ones. Therefore, for this particular specimen, we roughly estimate a characteristic energy scale in terms of the TLS Kondo model: Firstly, we consider  $T_S \simeq 0.2 \text{ K}$  as the temperature of the breakdown of the nFL scaling  $\Delta \rho \propto T^{1/2}$ . Further, we take  $T_{\rm K} = 5$  K, i.e, slightly more than an average value of  $T^* \simeq 4$  K inferred from Figs. 1 and 2. Note that  $T^*$  represents a low-temperature limit of the  $-\log T$  dependence that is expected to hold at  $T_{\rm K} < T$ . Finally, we get a crude estimate of the energy splitting  $\Delta_{\text{TLS}} \approx (T_{\text{K}} T_{\text{S}})^{1/2} \approx 1 \text{ K}$ . Thus, the  $\Delta_{\text{TLS}} \ll T_{\text{K}}$ requirement for the two-channel Kondo effect seems to be fulfilled in the ThAsSe sample discussed. Nevertheless, the large value of the Kondo temperature remains unclear, although the recent theoretical studies have shown that  $T_{\rm K}$  of the order of 0.5 K can be

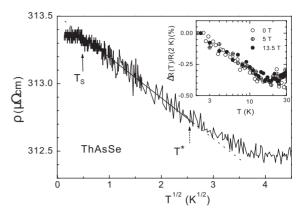


Fig. 2. The *ab*-plane resistivity of the structurally disordered ThAsSe showing a  $T^{1/2}$  behavior in a temperature window  $T_{\rm S} < T < T^*$ , as indicated by the solid line. Inset: The low-T upturn for another single crystal of ThAsSe shown on a semi-logarithmic temperature scale. The results, plotted as  $\Delta \rho(T)/\rho(2~{\rm K})$ , were obtained along the *ab*-plane in varying magnetic fields applied along the c-axis.

achieved outside the tunneling regime, i.e., when the first exited state is above the barrier [5].

To conclude, a complex low-T behavior of  $\rho(T)$  in the structurally disordered ThAsSe was observed. Though the enhanced energy scale, our  $\rho(T)$  results appear to be consistently interpreted in terms of the TLS Kondo model. In particular, indications towards a two-channel Kondo effect and the resultant nFL properties  $\Delta\rho \propto 1-aT^{1/2}$  derived from structural two-level systems were found for the ThAsSe single crystal with the very small value of  $T_{\rm S} \simeq 0.2$  K.

### References

- [1] D.L. Cox, A. Zawadowski, Adv. Phys. 47 (1998) 599.
- [2] T. Cichorek, et al., Phys. Rev. B 68 (2003) 144411.
- [3] T. Cichorek, et al., Phys. Stat. Sol. (b) 236 (2003) 351.
- [4] J. Schoenes, et al., Solid State Commun. 68 (1988) 287.
- [5] L. Borda, et al., cond-mat/0302334.