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# Suppression of the Kondo state in YbRh<sub>2</sub>Si<sub>2</sub> by large magnetic fields

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## Abstract

We present DC-magnetization,  $M(B)$ , magnetostriction,  $\Delta L(B)/L$ , and magnetoresistance  $\rho(T, B)$  measurements on high-quality single crystals of YbRh<sub>2</sub>Si<sub>2</sub> in magnetic fields up to 18 T and at temperatures down to 15 mK. At  $B^* \approx 9.5$  T, both  $M(B)$  and  $\Delta L(B)/L$  show pronounced changes of their slopes, indicative for a broadened phase transition. For fields above  $B^*$ , the coefficient  $A$  of the Fermi liquid behavior  $\Delta\rho = \rho_0 + A(B)T^2$  is reduced to very small values, indicating the suppression of the Heavy Fermion state.

High-field studies in Ce- and U-based compounds have revealed interesting phenomena such as metamagnetic transitions in CeRu<sub>2</sub>Si<sub>2</sub> and UPt<sub>3</sub> [1,2]. In the case of Yb-based valence-fluctuating compounds with characteristic temperatures  $T_0 \gtrsim 70$  K like YbCu<sub>5-x</sub>Ag<sub>x</sub>, the application of large magnetic fields induces a metamagnetic-like cross-over to a stable Yb<sup>3+</sup> state with localized magnetic moments [3]. Here we report on the high-field behavior of YbRh<sub>2</sub>Si<sub>2</sub> which is the only stoichiometric Yb-based HF system with a characteristic Kondo temperature of the order of 25 K [4]. It is located very close to a quantum critical point (QCP), related to very weak antiferromagnetic order at  $T_N = 70$  mK. A very small critical magnetic field  $B_c = 0.06$  T, applied perpendicular to the tetragonal  $c$ -axis, i.e. in the magnetic easy-plane, is sufficient to drive the system through the field-induced QCP [5]. For  $B > B_c$  a Landau Fermi liquid (LFL) state is induced with a strongly field-dependent heavy quasiparticle mass [5].

In order to study the high-field properties in YbRh<sub>2</sub>Si<sub>2</sub>, we performed DC-magnetization  $M(B)$ ,

magnetostriction  $\Delta L(B)/L$ , and magnetoresistance  $\rho(T, B)$  measurements on high-quality single crystals ( $\rho_0 = 1 \mu\Omega \text{ cm}$ ) of YbRh<sub>2</sub>Si<sub>2</sub>, prepared as described earlier [4]. For the magnetization and magnetostriction measurements a high-resolution Faraday magnetometer and a CuBe dilatometer have been adapted to dilution refrigerators, respectively. The resistivity was measured with the standard four-terminal AC technique.

The low- $T$  magnetization  $M(B)$  shows two anomalies (Fig. 1). The low-field anomaly near at 0.09 T is related to the QCP and indicates the polarization of a small moment of  $0.1\mu_B$ . The remaining moment is still fluctuating and contributes to the strongly enhanced Pauli-paramagnetic susceptibility [5]. Here, we concentrate on the high-field anomaly. The susceptibility  $\chi = dM/dB$  and magnetostriction coefficient  $\lambda = d(\Delta L(B)/L)/dB$  show step-like anomalies at  $B^* \approx 9.5$  T, indicative for a broadened second-order phase transition. The polarization at  $B^*$  amounts to  $1\mu_B$ . With increasing temperature, the kink in  $M(B)$  becomes broader but  $B^*$  is not shifted up to 2 K. For  $B^*$  the magnetostriction coefficient  $\lambda$  is negative, indicating a shrinking of the Yb-ions. Since the ionic radius of magnetic Yb<sup>3+</sup> configuration is smaller than that of the non-magnetic

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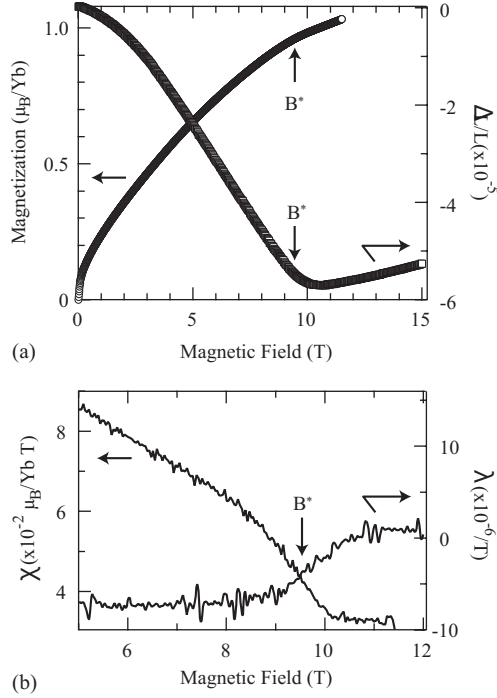


Fig. 1. Magnetization (at 80 mK, left axis) and magnetostriction (at 15 mK, right axis) vs magnetic field ( $B \perp c$ ) (a). Susceptibility  $\chi = dM/dB$  and magnetostriction coefficient  $\lambda = d(\Delta L(B)/L)/dB$  vs  $B$  (b). Arrows indicate critical field  $B^*$ .

$\text{Yb}^{2+}$  one, the effective valency increases with increasing field and reaches  $3+$  at  $B^*$ . The localization of the  $f$ -electrons leads to a reduction of the Pauli-paramagnetic contribution to the susceptibility and therefore results in a kink of the magnetization curve.

To get more information on the properties of the heavy quasiparticles around  $B^*$ , we analyze the field dependence of the coefficient  $A$ , derived from the LFL behavior of the electrical resistivity [5]. As shown in Fig. 2,  $A(B)$  is drastically reduced upon increasing  $B$  from below to above  $B^*$ . Since it has been shown that

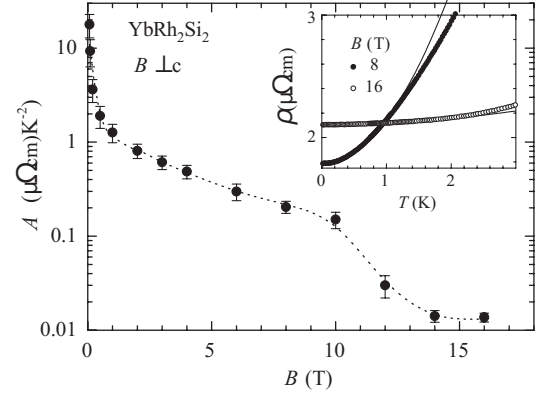


Fig. 2. Magnetic-field dependence of coefficient  $A$  (on a logarithmic scale) from Fermi liquid behavior  $\rho(T) = \rho_0 + AT^2$  in the electrical resistivity, observed for  $B > 0.06$  T ( $B \perp c$ ) [5]. Dotted line is intended as guide to the eyes. Inset shows  $\rho$  vs.  $T$  for 8 and 16 T. Solid lines represent  $T^2$  behavior.

the scaling relation  $A \propto \gamma^2$  holds at least up to 4 T [5], this indicates a step-like decrease of the quasiparticle mass. Using the value for  $A/\gamma^2$  determined in [5], a  $\gamma$ -coefficient of about  $70 \text{ mJ mol}^{-1} \text{ K}^{-2}$ , only, is estimated at 16 T.

To conclude, a broadened phase transition at  $B^* = (9.5 \pm 0.5)$  T is observed in  $\text{YbRh}_2\text{Si}_2$  for fields applied in the easy magnetic plane, which indicates the complete localization of the  $4f$ -electrons and the suppression of the HF state.

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