Two-Dimensional Superconducting Fluctuations in MgB$_2$ Films

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The origin of the resistive transition broadening for superconducting MgB$_2$ films is investigated experimentally. The crucial role of two-dimensional weak and critical fluctuations is demonstrated.

KEY WORDS: fluctuations; dimensionality; magnesium diboride.

1. INTRODUCTION

The discovery of superconductivity in MgB$_2$ [1], a material with hexagonal layered crystal structure, raised several questions about its transport properties. The crystal structure of MgB$_2$ and band structure calculations suggest that the quasi-two-dimensional boron planes are mainly responsible for the electrical transport. Therefore, the superconducting properties of magnesium diboride are expected to reflect a two-dimensional 2D character.

There exists a fundamental reason for the broadening of the resistive transition of a superconductor because of an intrinsic transition width associated with thermodynamic fluctuations of the order parameter. This minimal width, ΔT_c is given by the Ginzburg criterion [2], $\Delta T_c = Gi T_c$, where

$$Gi = \left[2\mu_0 \, k_B \, T_c / B_c^2(0)\xi_0^3 \right]^2 \tag{1}$$

is the Ginzburg parameter, with $B_c(0)$ the critical flux density at $T = 0$ K, $\xi_0 = \hbar v_F / 2\pi k_B T_c$, $\mu_0 = 4\pi \cdot 10^{-7}$ Vs/Am, and v_F the Fermy velocity. The value of Gi is extremely small for pure 3D conventional superconductors like bulk Pb and Al ($Gi \sim 10^{-13}$), rising many orders of magnitude for dirty and low dimensional systems making the fluctuation effects observable experimentally [3]. For layered superconductors with a small coherence length ξ_0, the intrinsic width ΔT_c may be much larger, up to $Gi \sim 10^{-2}$ for high-T_c superconductors like YBa$_2$Cu$_3$O$_{7-x}$ [4].

The layered character of the new superconductor MgB$_2$ with its small coherence length lets one expect the possibility of a detection of fluctuation effects. Therefore, in this work we report about an experimental study of the origin of the resistive transition broadening in MgB$_2$ thin films.

We performed critical field and resistance measurements of MgB$_2$ films grown on MgO single crystalline substrate by dc-magnetron sputtering. The linear temperature dependence of the upper critical field demonstrates a 3D character of superconductivity in our samples below T_c as discussed in [5]. In contrast, resistivity measurements yield a temperature dependence of the fluctuation conductivity above T_c which agrees with the Aslamazov–Larkin (AL) theory [6] for a 2D superconductor in the weak fluctuations region, $T > T_c$ and with the theory of critical fluctuations [7] in the vicinity of the superconducting transition, $T \sim T_c$ as shown in the present paper. Our experimental results indicate a quasi 2D origin of the nucleation of superconductivity in MgB$_2$.

2. SAMPLES

The MgB$_2$ films were prepared by dc-magnetron sputtering from Mg-MgB$_2$ composite targets on single

crystalline MgO substrate using a two-step process. The substrate temperature was kept at 200°C during the process of sputtering and then raised up to 600°C, followed by an *ex situ* annealing. Details of the sample preparation are given in [5]. The X-ray diffraction study yields a (101)-textured structure of the films, SEM-characterization of the samples demonstrated a polycrystalline very smooth morphology of the films with ~100 nm crystallites size. The resistive transitions $R(T)$ in constant external magnetic fields were measured by a conventional four-terminal resistive method. The special procedure used for determination of the residual resistance, R_n, is described in detail in [5].

3. RESULTS AND DISCUSSION

Figure 1 shows the resistive transition $R(T)$ at zero magnetic field of MgB$_2$ for one of the investigated samples. The transition width ΔT_c is about 0.7 K ($0.1R_n$–$0.9R_n$ criteron), demonstrating a high homogeneity of the film. In many works the broadening of the resistive transition of thin films and layered superconductors in zero magnetic field was interpreted in terms of superconducting fluctuations, rising in the vicinity of the critical temperature, T_c. The fluctuation or excess conductivity, $\sigma' \equiv 1/R(T) - 1/R_n$ strongly depends on the superconductor dimensionality, D.

For a 2D superconductor ($D = 2$) in the weak fluctuation region the excess conductivity $\sigma' \sim$
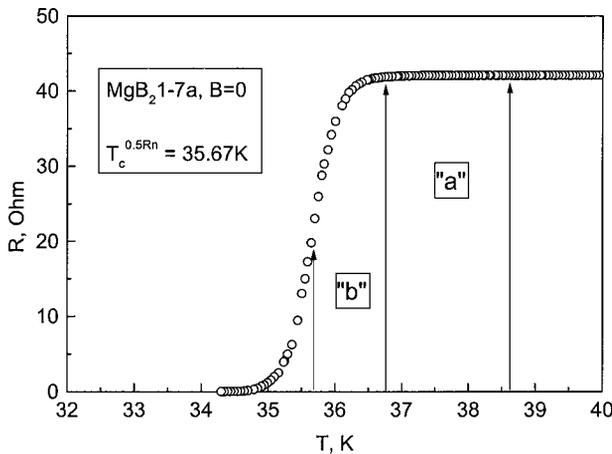
Fig. 1. Resistive transition for a MgB$_2$ film (thickness $d = 5.6\ \mu$m). Weak- and strong fluctuation regions are marked as "a" and "b," respectively.
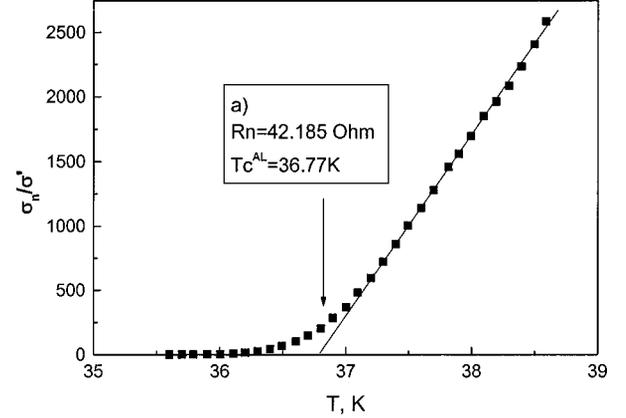
Fig. 2. Linear behavior of the fluctuation conductivity on 1/T in the weak fluctuation region, i.e. the upper part of the resistive transition, marked as "a" in Fig. 1. The solid line corresponds to the AL-theory for a 2D-superconductor, crossing the T-axis at T_c^{AL} (pointed by the arrow.)

$(T/T_c^{AL} - 1)^{(D-4)/2}$ is reverse proportional to the temperature

$$[\sigma'(T)]^{-1} = (R_n/\tau_{AL})(T - T_c^{AL})/T_c^{AL} \qquad (2)$$

where $\tau_{AL} = (R_n^\square e^2)/16\hbar$ and $R_n^\square$ is the sheet resistance of the film [6].

Figure 2 presents the results of Fig. 1 in the (σ_n/σ') on T plot, indicating that the resistive behavior of the MgB$_2$ film is caused by superconducting fluctuations according to Eq. (2). The fit of the experimental results in Fig. 2 by Eq. (2) gave us the values of $R_n^\square$ and T_c^{AL}. The intersection of the linear fit, according to Eq. (2), with the T-axis gives the critical temperature $T_c^{AL} = 36.77$ K, higher than the middle-point value of $T_c^{0.5R_n} = 35.67$ K as a result of the critical temperature shift due to fluctuations.

Figure 3 shows the results of Fig. 1 in the $ln(\sigma_n/\sigma')$ on T plot. According to the theory of critical fluctuations [7] in the vicinity of T_c for a 2D superconductor an exponential behavior of the fluctuation conductivity is expected, with

$$[\sigma'(T)]^{-1} \sim R_n \exp[-(1 - T/T_c)/Gi] \qquad (3)$$

The solid line in Fig. 3 is the best fit of the experimental data to Eq. (3). From the inset of Fig. 3, we get the slope $B = 1/Gi = 76.3774$ yielding the value $Gi = 0.013$ for the Ginzburg parameter. This value results in the resistive transition width $\Delta T_c = GiT_c = 0.47$ K which is in a good agreement with the experimental value of $\Delta T_c = 0.7$ K, mentioned above.
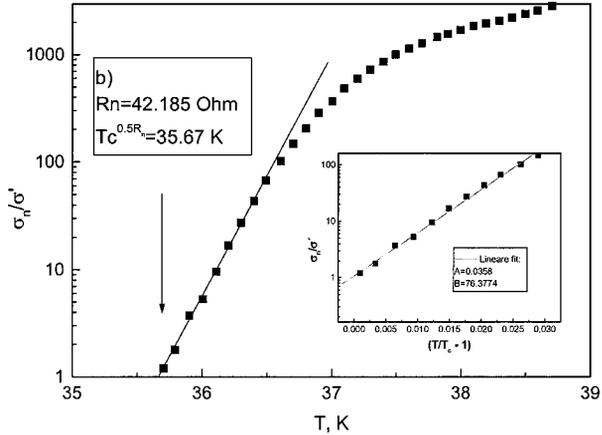
Fig. 3. Linear behavior of the fluctuation conductivity in a semilogarithmic scale on temperature in the strong fluctuation region, marked as "b" in Fig. 1.

4. CONCLUSIONS

From our experimental results follows a strong influence of superconducting fluctuations on the broadening of the resistive transition of superconducting MgB_2. The 2D character of the fluctuations as well in the weak fluctuation region (Aslamazov–Larkin case) as also in the critical fluctuation region demonstrates the two-dimensional nucleation of superconductivity in this novel layered compound.

REFERENCES

1. J. Namagatsu, N. Nagakawa, T. Muranaka, Y. Zenitani, and J. Akimitsu, *Nature* **410,** 63 (2001).
2. V. V. Schmidt, *Introduction to the Physics of Superconductors*, 2nd Russian edn., enlarged by V. V. Ryazanov and M. F. Feigelman (MCNMO Publishers, Moskow, 2000) Chapter III, §19.
3. N. Ya. Fogel, A. S. Sidorenko, L. F. Rybalchenko, and I. M. Dmitrenko, *Sov. Phys. JETP* **50,** 120 (1979).
4. A. Kapitulnik, *Physica C* **153–155,** 520 (1988).
5. A. S. Sidorenko, L. R. Tagirov, A. N. Rossolenko, N. S. Sidorov, V. I. Zdravkov, V. V. Ryazanov, M. Klemm, S. Horn, and R. Tidecks, *JETP Lett.* **76,** 17 (2002).
6. L. G. Aslamazov and A. I. Larkin, *Phys. Lett.* **26A,** 238 (1968).
7. L. P. Kadanoff and G. Laramore, *Phys. Rev.* **175,** 579 (1968).