

Microscopic origin of the linear temperature increase of the magnetic susceptibility of BaFe₂As₂S. L. Skornyakov,^{1,2} V. I. Anisimov,^{1,2} and D. Vollhardt³¹*Institute of Metal Physics, Russian Academy of Sciences, 620041 Yekaterinburg GSP-170, Russia*²*Ural Federal University, 620002 Yekaterinburg, Russia*³*Theoretical Physics III, Center for Electronic Correlations and Magnetism, Institute of Physics, University of Augsburg, D-86135 Augsburg, Germany*

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Employing a combination of *ab initio* band-structure theory and dynamical mean-field theory we explain the experimentally observed linear temperature increase of the magnetic susceptibility of the iron pnictide material BaFe₂As₂. The microscopic origin of this anomalous behavior is traced to a sharp peak in the spectral function located approximately 100 meV below the Fermi level. This peak is due to the weak dispersion of two-dimensional bands associated with the layered crystal structure of pnictides.

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I. INTRODUCTION

Since the discovery of high-temperature superconductivity in the iron pnictides¹ the unusual electronic properties of this class of materials has attracted considerable attention.^{2–13} The interest was stimulated by the fact that the new superconductors share several similarities with the well-studied, but still not well-understood, high- T_c cuprates. First, both classes of superconductors crystallize into a layered structure. Second, in most cases the parent compounds of the pnictides are not superconducting, and superconductivity emerges only under doping or pressure and is associated with the suppression of antiferromagnetic (AFM) order. However, in contrast to the cuprates whose parent compounds are Mott insulators, the parent compounds of the pnictides are multi-band metals. The nature of the magnetic ground state of the parent compounds is also different: In the cuprates it corresponds to a Néel-type order of a Mott-Hubbard insulator, while in the pnictides magnetism is associated with a nesting-induced spin density wave^{2,3,14} (SDW).

The magnetic properties of pnictide materials show anomalous behavior even in the paramagnetic state. An unusual linear temperature increase of the uniform magnetic susceptibility was reported in the parent compound BaFe₂As₂ (Ref. 4) as well as in stoichiometric and fluorine-doped LaFeAsO (Ref. 5). It is now well established that the linear increase of the magnetic susceptibility with temperature is a general property of all pnictide superconductors for temperatures above the SDW transition. Nevertheless, no consensus has been reached so far about the origin of this phenomenon. To date several mechanisms were proposed to explain the observed T -dependence in the pnictides. Wang *et al.*⁴ and Zhang *et al.*⁶ suggested that the linear- T behavior is a consequence of strong antiferromagnetic fluctuations present above the SDW transition temperature. Korshunov *et al.*⁷ argued that short-range antiferromagnetic fluctuations are the source of a linear- T term in the susceptibility of a two-dimensional Fermi liquid which allowed them to obtain good agreement with the experimental data.

A very important issue concerning the spectral and magnetic properties of the pnictides is the role of Coulomb correlations. It is now generally accepted that electronic correlations in the pnictides are not as strong as in the cuprates and should

be classified as moderate.^{8,9} It was shown^{9,10} that the spectral properties of the pnictides can be reproduced by first-principles techniques only if local dynamical Coulomb correlations are taken into account. This can be achieved by employing the LDA + DMFT method.¹⁵ This computational scheme combines electronic band-structure calculations in the local density approximation (LDA) with the many-body physics incorporated in the dynamical mean-field theory (DMFT).¹⁶

In our earlier study¹¹ we proposed an explanation of the temperature increase of the magnetic susceptibility of LaFeAsO which was based on a first-principles analysis of the low-energy spectral properties caused by local dynamical Coulomb correlations, without taking into account interatomic magnetic fluctuations. In the present paper we employ the LDA + DMFT scheme to demonstrate that the proposed mechanism can be applied also to understand the origin of the linear temperature dependence of the uniform magnetic susceptibility in stoichiometric BaFe₂As₂.

II. COMPUTATIONAL METHOD

The LDA + DMFT computational scheme implemented in the present work proceeds in four steps: (i) the construction of an effective tight-binding Hamiltonian $\hat{H}^{\text{WF}}(\mathbf{k})$ from a converged LDA solution by projecting¹⁷ onto Wannier functions, (ii) the addition of the local Coulomb interaction \hat{H}^{Coul} , (iii) a double-counting correction which takes into account the local interactions already described by the LDA, and (iv) the self-consistent solution of the DMFT equations on the Matsubara contour with continuous-time quantum Monte Carlo¹⁸ (CT-QMC) as impurity solver. The Hamiltonian to be solved by DMFT is given by

$$\hat{H}^{\text{DMFT}}(\mathbf{k}) = \hat{H}^{\text{WF}}(\mathbf{k}) + \hat{H}^{\text{Coul}} + \hat{H}^{\text{DC}}. \quad (1)$$

Exact investigations of the magnetic response of this model must employ a rotationally invariant form of the interaction term \hat{H}^{Coul} . However, up to now there did not exist effective algorithms for the solution of a five-orbital Hubbard model with the full Coulomb interaction within CT-QMC. Furthermore, our investigation of the magnetic properties of BaFe₂As₂ requires a separate, and extremely time-consuming, *self-consistent* DMFT calculation for each point of the

susceptibility curve. To make computations feasible, we therefore include only the density-density contributions to the full interaction

$$\hat{H}^{\text{Coul}} \equiv \hat{H}^U = \frac{1}{2} \sum_{i,\alpha,\alpha',\sigma,\sigma'} U_{\alpha\alpha'}^{\sigma\sigma'} \hat{n}_{i\alpha\sigma}^d \hat{n}_{i\alpha'\sigma'}^d. \quad (2)$$

Here $U_{\alpha\alpha'}^{\sigma\sigma'}$ is the Coulomb interaction matrix, $\hat{n}_{i\alpha\sigma}^d$ is the occupation number operator for d electrons in the orbital α or α' , with spin σ or σ' , on the i th site. This approximation neglects spin flip and pair hopping processes. Nevertheless, as will be shown below, it is able to provide correct results for the spectral and magnetic properties of BaFe_2As_2 . The double-counting term is $\hat{H}^{\text{DC}} = -\bar{U}(n_{\text{DMFT}} - \frac{1}{2})\hat{I}$, where n_{DMFT} is the total, self-consistently determined number of d electrons obtained within LDA + DMFT, and \bar{U} is the average Coulomb parameter for the d shell.

We construct Wannier functions in the energy window including Fe- d and As- p states. Hence, by construction energy bands of the H^{WF} Hamiltonian exactly reproduce 16 Fe- d and As- p bands (two As and Fe atoms in the formula unit, one formula unit in the unit cell) obtained in LDA calculations, and the p - d hybridization is explicitly taken into account.

The interaction matrix $U_{\alpha\alpha'}^{\sigma\sigma'}$ is parametrized by the effective on-site Coulomb parameter U and intra-orbital exchange parameter J according to the procedure described in Ref. 19. In the present calculation we use $U = 3.5$ eV and $J = 0.85$ eV obtained with the constrained DFT procedure.^{20–22}

The orbitally resolved spectral functions of the interacting system are then computed as

$$A_\alpha(\omega) = -\frac{1}{\pi} \text{Im} \sum_{\mathbf{k}} [(\omega + \mu)\hat{I} - \hat{H}^{\text{WF}}(\mathbf{k}) - \hat{H}^{\text{DC}} - \hat{\Sigma}(\omega)]_{\alpha\alpha}^{-1}, \quad (3)$$

where the subscript α refers to an orbital, μ is the self-consistent chemical potential, and $\Sigma(\omega)$ is the self-energy on the real axis obtained by analytic continuation using the Padé approximant²³ technique; the details are described in Ref. 10.

The uniform magnetic susceptibility is calculated as the response of the system to a weak external magnetic field,

$$\chi(T) = \frac{\partial M(T)}{\partial E_h} = \frac{\partial [n_\uparrow(T) - n_\downarrow(T)]}{\partial E_h}, \quad (4)$$

where $M(T)$ is the field-induced magnetization, $n_\sigma(T)$ is the number of electrons with spin σ , and E_h is the energy correction corresponding to the applied field. Since the field is finite the calculations are performed in three steps: First, we check that the polarization is zero in the absence of the field, then we check that $M(T)$ is a linear function of E_h , and finally we evaluate the derivative in Eq. (4) as a ratio of $M(T)$ and E_h .

III. RESULTS

A. Temperature dependence of the uniform magnetic susceptibility

In Fig. 1 the uniform magnetic susceptibility $\chi(T)$ computed within LDA + DMFT is compared with the experimental data of Wang *et al.*⁴ In the temperature range from 200 to 600 K the temperature dependence of the calculated $\chi(T)$ is found to be almost perfectly linear. However, the slope is by

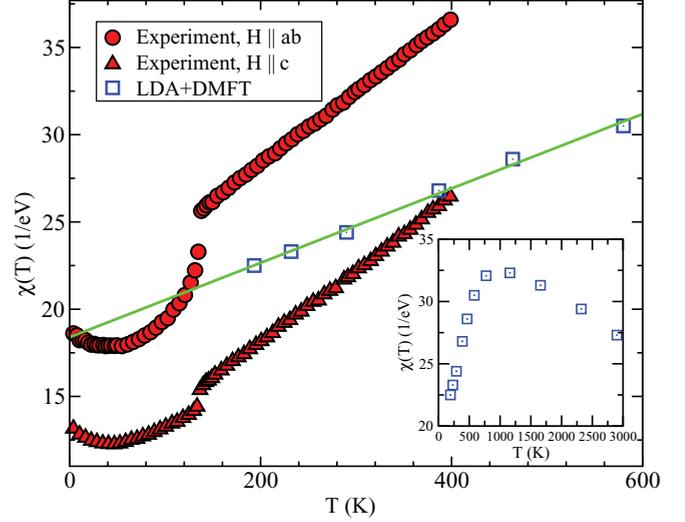


FIG. 1. (Color online) Uniform magnetic susceptibility $\chi(T)$ of BaFe_2As_2 calculated within LDA + DMFT (squares) in comparison with experimental data of Wang *et al.* (Ref. 4) (circles and triangles). The line is a least-square fit to the last six points of the computed data. The inset shows the theoretical curve for the full temperature interval.

a factor of 1.7 smaller than in the experiment. The origin of this quantitative discrepancy is not clear at the moment and will be the subject of future investigations. To emphasize the linearity we plot a least-square fit to the last six points of the computed data. We note that the observed linear behavior is, in fact, due to an extended linear region around the turning point of $\chi(T)$ at ≈ 350 K. The obtained $\chi(T)$ has a maximum at about 1000 K and decreases for higher temperatures.

More detailed information about the magnetic properties of BaFe_2As_2 can be obtained from an analysis of the orbitally resolved contributions $\chi_\alpha(T)$ ($\alpha = xy, yz, xz, 3z^2 - r^2, x^2 - y^2$) to the total magnetic susceptibility. In Fig. 2 we show the temperature dependence of the Fe 3d susceptibilities of BaFe_2As_2 obtained in LDA + DMFT. All contributions have

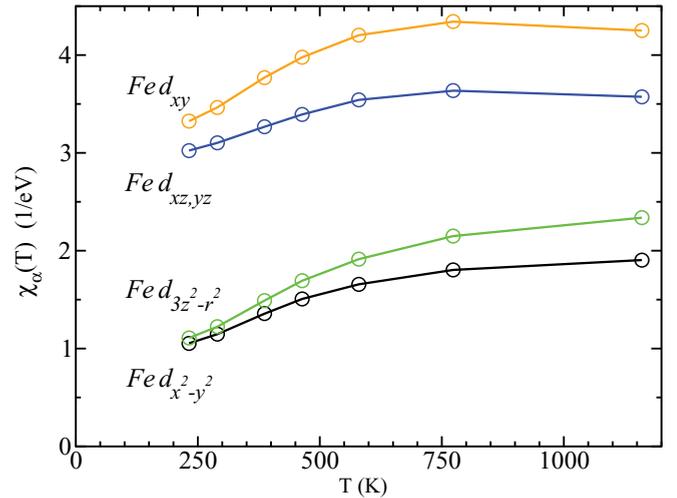


FIG. 2. (Color online) LDA + DMFT results for the orbitally resolved Fe 3d susceptibilities $\chi_\alpha(T)$ with $\alpha = xy, yz, xz, 3z^2 - r^2, x^2 - y^2$ of BaFe_2As_2 vs. temperature obtained from the derivative of the magnetization.

approximately equal slope in the temperature interval from 200 to 500 K. The d_{xy} orbital of Fe provides the largest contribution to the total susceptibility.

B. Connection between magnetic and spectral properties

In our previous paper¹¹ we proposed a scenario according to which the anomalous T behavior of $\chi(T)$ in the pnictides is connected with the presence of a sharp peak in the spectral function below the Fermi energy. In Fig. 3 the total Fe 3d spectral function $A(\omega)$ computed within LDA + DMFT is shown in comparison with the LDA result. It demonstrates that dynamical correlation effects strongly renormalize the spectrum in the vicinity of the Fermi energy. In particular, the electronic correlations are seen to lead to a narrow peak below the Fermi level while the remaining part of the spectrum is only weakly affected by the correlations. The peak in the energy window from -4 to -2 eV should not be mistaken for a lower Hubbard band since a similar peak is already present in the LDA result. The Fe 3d spectral weight in this energy area is a consequence of the hybridization with As 4p states.

A comparison of the orbitally resolved spectral functions computed within LDA and LDA + DMFT, respectively, is shown in Fig. 4. The LDA + DMFT results are in good agreement with previously reported theoretical and experimental spectra.⁹ Except for the $d_{x^2-y^2}$ orbital the spectral functions obtained by LDA + DMFT all show a sharp peak below the Fermi energy. These peaks originate from local correlation effects as pointed out above.

A quantitative measure of the correlation strength is the quasiparticle renormalization factor Z . In the single-orbital case it can be obtained from the real-axis self-energy $\Sigma(\omega)$ which is related to the effective mass enhancement m^*/m by $Z^{-1} = 1 - \partial \text{Re}(\Sigma(\omega))/\partial \omega = m^*/m$. For a multi-orbital problem the self-energy is a matrix. Therefore different orbitals have different m^*/m . The computed values of the mass enhancement range from 2.5 to 3.7 and agree well with previous estimates of m^*/m for pnictides obtained from the renormalization of the LDA band structure.^{8-10,12,13}

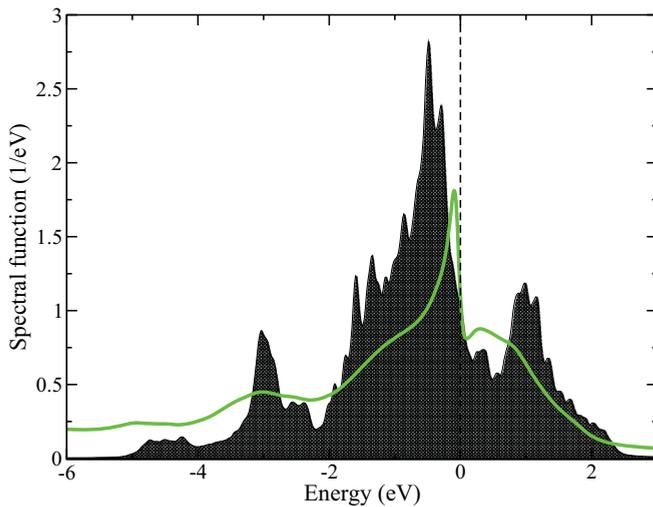


FIG. 3. (Color online) Comparison of the total Fe 3d spectral function of BaFe_2As_2 as obtained from LDA + DMFT (solid line) and LDA (shaded area), respectively.

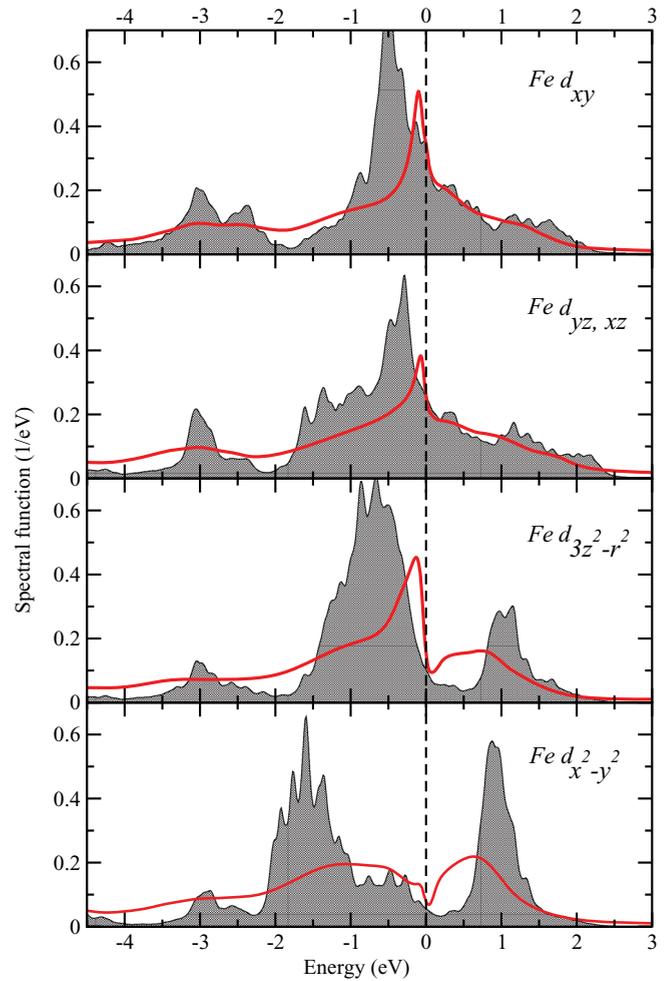


FIG. 4. (Color online) Orbitally resolved Fe 3d spectral functions of BaFe_2As_2 obtained within LDA + DMFT (solid lines) in comparison with LDA results (filled areas). The Fermi energy is set to 0 eV.

To identify possible reasons for the anomalous behavior of the magnetic properties it is instructive to compute the temperature evolution of the orbitally resolved spectral functions. The spectral curves computed for temperatures ranging from 232 to 580 K are shown in Fig. 5. All spectral functions, except for the $d_{x^2-y^2}$ orbital, show a temperature sensitive peak which is located approximately 100 meV below the Fermi energy. These peaks increase in amplitude and become narrow with decreasing T . We note that the magnitudes of the orbital contribution $\chi_\alpha(T)$ are proportional to the density of states of the corresponding orbitals at the Fermi energy.

The simplest way to establish a possible connection between the temperature evolution of the magnetic susceptibility and the excitation spectrum in the multi-orbital case is to estimate the susceptibility (per spin) using the bubble diagram obtained by convoluting the DMFT Green's functions, Eq. (5),

$$\chi_\alpha^0(T) = \frac{1}{\beta} \sum_{\mathbf{k}, i\omega, \alpha'} \hat{G}_{\alpha\alpha'}(\mathbf{k}, i\omega) \hat{G}_{\alpha'\alpha}(\mathbf{k}, i\omega), \quad (5)$$

where $\hat{G}_{\alpha\alpha'}(\mathbf{k}, i\omega) = [(i\omega + \mu)\hat{I} - \hat{H}^{\text{DMFT}}(\mathbf{k}) - \hat{\Sigma}(i\omega)]_{\alpha\alpha'}^{-1}$. This expression describes the spin susceptibility in the absence of vertex corrections (i.e., gives an estimate of the

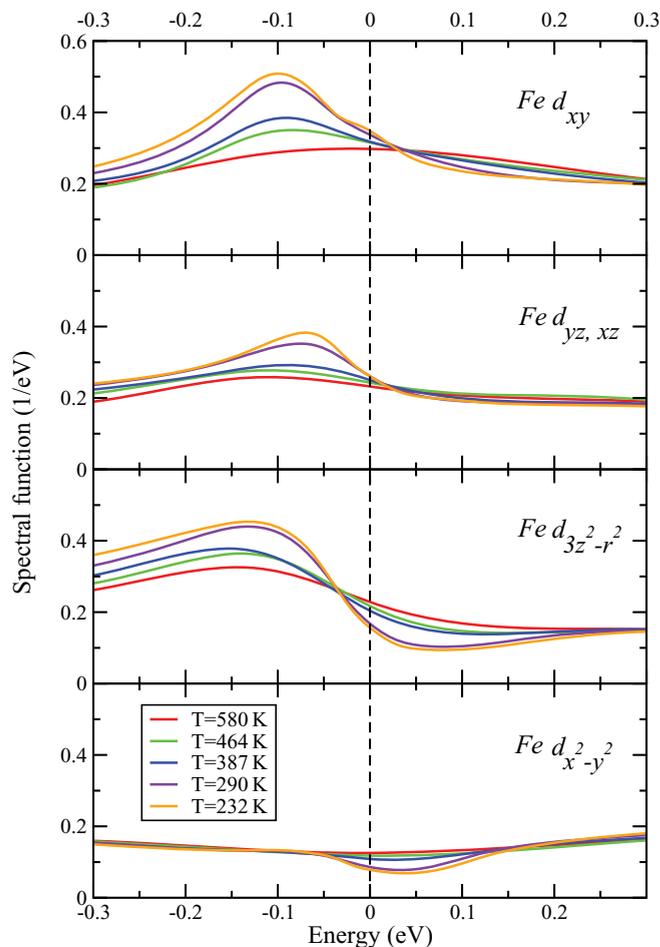


FIG. 5. (Color online) Fe 3d spectral functions of BaFe₂As₂ computed within LDA + DMFT in the temperature range 232–580 K. The Fermi energy corresponds to 0 eV.

magnetic response due to single-particle excitations which are characterized by the interacting spectral function). An explicit connection between the magnetic response, Eq. (5), with the excitation spectrum can be made in the one-orbital case when off-diagonal elements of \hat{G} are absent:

$$\chi_1(T) = -\frac{1}{4\pi^2} \int d\omega_1 d\omega_2 \frac{f_F(\omega_1) - f_F(\omega_2)}{\omega_1 - \omega_2} A(\omega_1)A(\omega_2). \quad (6)$$

Here $A(\omega)$ is the spectral function of the interacting system (per spin), and the temperature enters via the Fermi function, $f_F(\omega)$. In real compounds multi-orbital physics, including hybridization effects, is important, in which case the full matrix Green functions must be used in Eq. (5).

The influence of interaction effects may be estimated by calculating the magnetic susceptibility in the random-phase approximation (RPA). If the orbital dependence of the Coulomb interaction between d electrons is neglected, the orbital contributions to the total uniform d magnetic susceptibility within RPA are given by

$$\chi_\alpha^{\text{RPA}}(T) = \frac{2\chi_\alpha^0(T)}{1 - \bar{U}\chi_\alpha^0(T)}, \quad (7)$$

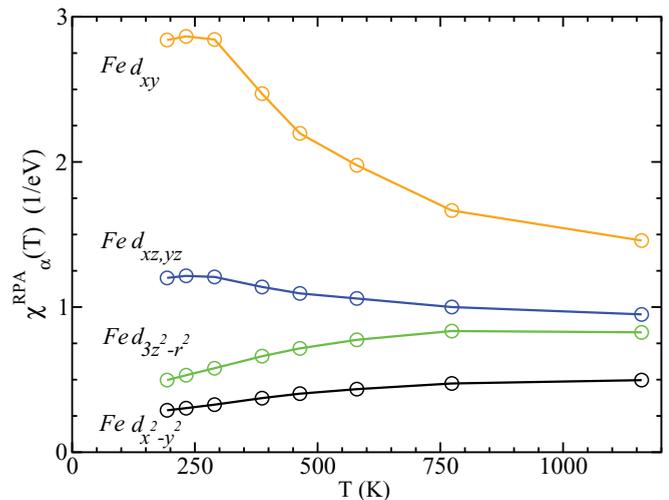


FIG. 6. (Color online) LDA + DMFT results for the orbitally resolved Fe 3d susceptibilities $\chi_\alpha^{\text{RPA}}(T)$ with $\alpha = xy, yz, xz, 3z^2 - r^2, x^2 - y^2$ of BaFe₂As₂ vs. temperature, obtained from the RPA expression for the susceptibility in Eq. (7).

where the factor 2 is due to the spin degeneracy. The temperature dependence of the orbitally resolved susceptibilities computed according to Eq. (7), is shown in Fig. 6. For the $d_{3z^2-r^2}$ and $d_{x^2-y^2}$ orbitals the results are in good qualitative agreement with the full LDA + DMFT solution. By contrast, the temperature dependence of the susceptibilities corresponding to the d_{xy} , d_{yz} , and d_{xz} states is not reproduced by Eq. (7); Apparently vertex corrections are important in this case.

The above results suggest the following interpretation, whose correctness will be demonstrated later: The electronic states forming the sharp peak in the spectral function below the Fermi energy lead to thermal excitations which contribute to the susceptibility. When the energy $k_B T$ is larger than the distance between the peak and the Fermi level, the number of states which can be excited is reduced and the susceptibility starts to decrease. A more complex mechanism is responsible for the increase of the $d_{x^2-y^2}$ susceptibility where the corresponding spectral function does not show a peak below the Fermi energy.

Namely, as in the case of LaFeAsO (Ref. 11) the off-diagonal contributions to the susceptibility, $\chi_{\alpha\alpha'}^0 = 1/\beta \sum_{\mathbf{k}, i\omega} \hat{G}_{\alpha\alpha'}(\mathbf{k}, i\omega) \hat{G}_{\alpha'\alpha}(\mathbf{k}, i\omega)$, and in particular $\chi_{x^2-y^2, 3z^2-r^2}^0(T)$, are responsible for the increase of $\chi_{x^2-y^2}^0(T)$. Thus the temperature increase of $\chi_{x^2-y^2}^{\text{RPA}}(T)$ is caused by the magnetic response of the other orbitals (for details see the supplementary material of Ref. 11). A detailed analysis of the proposed mechanism of the increase of $\chi(T)$ using a simplified model will be presented in the following section.

C. Model analysis

To further clarify the relation between the shape of the spectral function and the anomalous temperature behavior of the magnetic properties we will now perform a model calculation, where multi-orbital effects are neglected. As a first step we compute the DMFT spin susceptibility for a single-

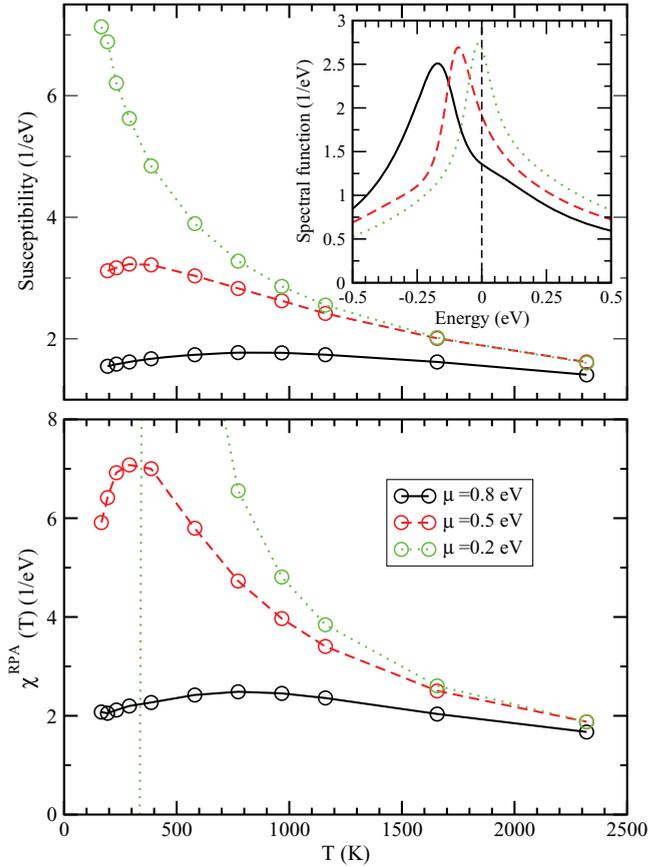


FIG. 7. (Color online) LDA + DMFT results for the uniform magnetic susceptibility of a one-band model with the noninteracting Fe d_{xy} DOS of BaFe₂As₂ as a function of temperature. Upper panel: Results computed according to Eq. (4); the inset shows the spectral function of the interacting system. Lower panel: RPA results for the uniform magnetic susceptibility using Eq. (7); the divergence of the susceptibility for $\mu = 0.2$ eV at about 300 K is an artifact of the RPA.

band model, constructed in such a way that the noninteracting system has exactly the same spectral function as the one for the d_{xy} orbital of the tight-binding Hamiltonian H_{WF} in Eq. (1). In the upper panel of Fig. 7 the temperature dependence of the spin susceptibility computed for the model within DMFT is shown for several values of the chemical potential. The corresponding spectral functions of the interacting system are presented in the inset of Fig. 7. Depending on the peak position there are two characteristically different temperature dependencies of the susceptibility: (i) An increase in the low temperature region with a maximum at an intermediate temperature followed by a decrease at higher T ($\mu = 0.5, 0.8$ eV), and (ii) a monotonic decrease with temperature ($\mu = 0.2$ eV). The former regime is obtained when the peak is substantially below the Fermi energy, the latter regime corresponds to the case when the peak is very close to, or right at, the Fermi energy. The result of the estimate of the susceptibility according to Eq. (7) is presented in the lower panel of Fig. 7. It is seen to reproduce all features of the curves computed by Eq. (4).

This confirms that the magnetic response of the system is indeed governed by its spectral properties. Thus, the driving force of the nonmonotonic behavior of the susceptibility is

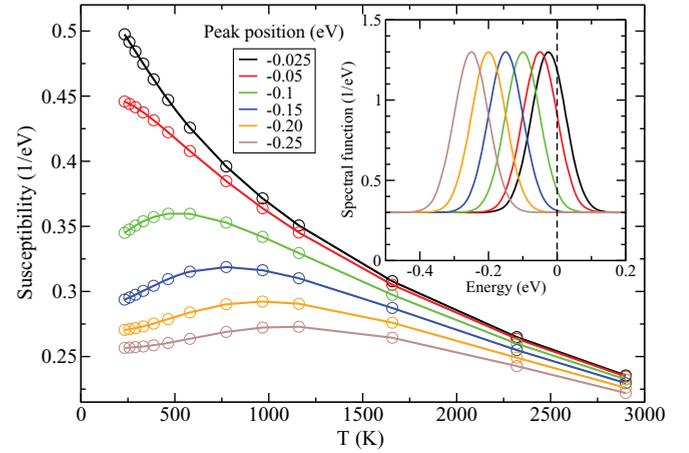


FIG. 8. (Color online) Spin susceptibility $\chi_1(T)$ of a model defined by the density of states (8) calculated from Eq. (6) for different values of the peak position. The inset shows the density of states plotted relative to the Fermi energy (zero energy).

the peak below the Fermi energy. In particular, the distance between the peak and the Fermi level is important. In spite of the fact that the peak in the spectral function exists at the level of LDA, its position is too far from the Fermi energy to cause a serious increase of the susceptibility. The correlations shift the peak towards the Fermi level which results in more pronounced temperature dependence of $\chi(T)$.

To relate the shape of the spectral function with the two temperature regimes of the magnetic susceptibility even more explicitly, we now compute the temperature behavior of $\chi^0(T)$ for a system with a density of states given by a Gauss function with offset from the energy axis

$$A(\omega) = A + B \exp \left[\left(\frac{\omega - C}{2\sigma} \right)^2 \right]. \quad (8)$$

The parameters in Eq. (8) were adjusted such that the maximum and the width of the function are close to the ones obtained with LDA + DMFT for the material specific Hamiltonian for BaFe₂As₂. The results for $\chi_1(T)$ computed according to Eq. (6) are shown in Fig. 8. The temperature behavior of $\chi_1(T)$ and its evolution upon changes of the peak position in the spectral function are seen to qualitatively reproduce all features obtained in DMFT for the one-band model. This can be viewed as a direct indication that the peculiarities observed in the anomalous behavior of the spin susceptibility in DMFT originate from the shape of the spectral function in the vicinity of the Fermi energy.

Finally we address the question concerning the microscopic origin of the peaks in the spectral function below the Fermi level. In the following analysis we focus on the d_{xy} orbital since the peak in its spectral function is sharper than those of the other orbitals. The contributions of the d_{xy} states to the band structure are shown in the left upper panel of Fig. 9 as “fat bands” (i.e., the thickness of a band is proportional to the contribution of states with selected symmetry). The peak in the d_{xy} spectral function (Fig. 9, right upper panel) is formed by regions of relatively flat bands centered at the energy approximately -0.4 eV relative to the Fermi level.

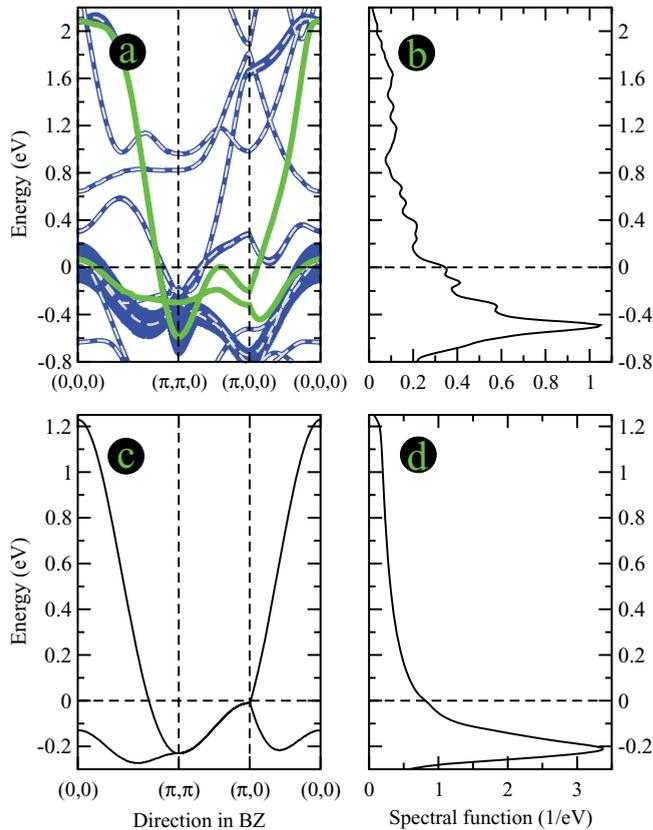


FIG. 9. (Color online) Band structure and spectral functions computed for BaFe_2As_2 ; the Fermi energy corresponds to 0 eV. (a) Dispersion curves calculated within LDA (dashed curves), contributions of the orbitals with d_{xy} symmetry (fat bands), energy bands of a two-orbital model obtained as a projection onto the d_{xy} states (solid curves). (b) Spectral function of the d_{xy} orbital from LDA. (c) Energy bands of a model corresponding to the real hopping parameters. (d) Spectral function of the model Hamiltonian (9).

To construct a minimal model describing the energy dispersion of the d_{xy} states we solve an effective two-band Hamiltonian $H^{2D}(\mathbf{k})$ with two Fe atoms in the unit cell, which is obtained from a projection of Bloch states in the vicinity of the Fermi energy onto a subspace of Wannier functions with d_{xy} symmetry. The energy bands of $H^{2D}(\mathbf{k})$ are shown in the left upper panel of Fig. 9 by solid curves. In the next step we introduce the real-space Hamiltonian H_{real}^{2D} written for a square lattice with two atoms in the unit cell and nonzero hoppings within three coordinate spheres. The Hamiltonian H_{real}^{2D} has the form

$$H_{\text{real}}^{2D} = t \sum_{i\mathbf{R}} c_{i\mathbf{R}}^\dagger c_i + t' \sum_{i\mathbf{R}'} c_{i\mathbf{R}'}^\dagger c_i + t'' \sum_{i\mathbf{R}''} c_{i\mathbf{R}''}^\dagger c_i, \quad (9)$$

where i labels the atoms and the radius vectors \mathbf{R} , \mathbf{R}' , and \mathbf{R}'' correspond to the cluster of nearest, next-nearest, and next-next-nearest neighbors, respectively, around atom i . The hopping parameters $t = -170$ meV, $t' = 98$ meV, and $t'' = 21$ meV were computed as Fourier transforms of the material-specific Hamiltonian $H^{2D}(\mathbf{k})$.

The shape of the energy bands and the spectral function computed for Hamiltonian (9) are shown in the lower panels

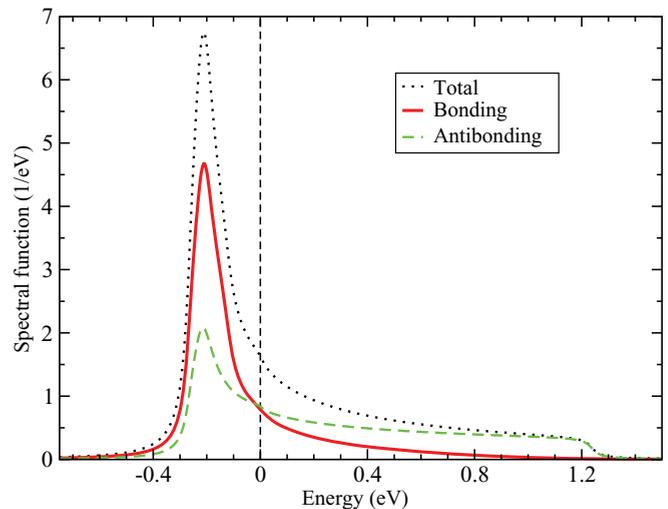


FIG. 10. (Color online) Spectral functions of Hamiltonian (9) calculated in the basis of bonding (solid curve) and antibonding (dashed curve) wave functions. The total spectral function is shown by a dotted line. The Fermi energy is set to zero.

of Fig. 9. They are in good agreement with the corresponding characteristics obtained from the direct calculation. As in the multi-orbital case, the band structure of Eq. (9) represents a combination of dispersive bands and bands with less pronounced \mathbf{k} dependence. In particular, it follows that the peak below the Fermi energy is formed by a relatively flat band located in the same energy interval as the peak.

To understand the coexistence of the flat and dispersive regions within a band of given symmetry (d_{xy} in our case) it is instructive to plot the spectral functions of the Hamiltonian (9) in the basis of bonding and antibonding wave functions. In other words, if $|\phi_1\rangle$ is the wave function on the one atom in the unit cell, $|\phi_2\rangle$ is the wave function of the other atom. Then the new basis is defined as $|\phi_+\rangle = |\phi_1\rangle + |\phi_2\rangle$ and $|\phi_-\rangle = |\phi_1\rangle - |\phi_2\rangle$. Spectral functions of the Hamiltonian (9) computed for that basis with realistic hopping parameters and chemical potential are shown in Fig. 10. From Fig. 10 it follows that the band with a pronounced dispersion is mainly formed by antibonding linear combination $|\phi_-\rangle$ and the main contribution to the peak is provided by the bonding function $|\phi_+\rangle$ whose dispersion is less pronounced.

IV. CONCLUSION

By employing the LDA + DMFT method we investigated the interplay between the spectral and magnetic properties of BaFe_2As_2 . The calculated temperature dependence of the uniform magnetic susceptibility is in good agreement with experimental data. Our calculations show that there are pronounced, temperature sensitive peaks below the Fermi energy in the spectral function of BaFe_2As_2 . We proposed a scenario according to which the temperature increase of the susceptibility is a consequence of the thermal excitation of the electronic states which lead to these peaks. Our analysis is based on the DMFT solution of a one-band model with a density of states corresponding to the Fe- d_{xy} spectral function

of BaFe_2As_2 . The results clearly demonstrate that the peak in the spectral function in the vicinity of the Fermi energy is a prerequisite for the linear temperature increase of magnetic susceptibility. The peaks in the real compound are due to the weak dispersion of bonding states arising from the layered structure of BaFe_2As_2 .

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- ¹Y. Kamihara, H. Hiramatsu, M. Hirano, R. Kawamura, H. Yanagi, T. Kamiya, and H. Hosono, *J. Am. Chem. Soc.* **128**, 10012 (2006); Y. Kamihara, T. Watanabe, M. Hirano, and H. Hosono, *ibid.* **130**, 3296 (2008); Z.-A. Ren, W. Lu, J. Yang, W. Yi, X.-L. Shen, Z.-C. Li, G.-C. Che, X.-L. Dong, L.-L. Sun, F. Zhou, and Z.-X. Zhao, *Chin. Phys. Lett.* **25**, 2215 (2008).
- ²D. J. Singh, *Physica C* **469**, 418 (2009).
- ³M. Rotter, M. Tegel, D. Johrendt, I. Schellenberg, W. Hermes, and R. Pottgen, *Phys. Rev. B* **78**, 020503 (2008).
- ⁴X. F. Wang, T. Wu, G. Wu, H. Chen, Y. L. Xie, J. J. Ying, Y. J. Yan, R. H. Liu, and X. H. Chen, *Phys. Rev. Lett.* **102**, 117005 (2009).
- ⁵R. Klingeler, N. Leps, I. Hellmann, A. Popa, U. Stockert, C. Hess, V. Kataev, H.-J. Grafe, F. Hammerath, G. Lang, S. Wurmehl, G. Behr, L. Harnagea, S. Singh, and B. Büchner, *Phys. Rev. B* **81**, 024506 (2010).
- ⁶G. M. Zhang, Y. H. Su, Z. Y. Lu, Z. Y. Weng, D. H. Lee, and T. Xiang, *Europhys. Lett.* **86**, 37006 (2009).
- ⁷M. M. Korshunov, I. Eremin, D. V. Efremov, D. L. Maslov, and A. V. Chubukov, *Phys. Rev. Lett.* **102**, 236403 (2009).
- ⁸M. Aichhorn, L. Pourovskii, V. Vildosola, M. Ferrero, O. Parcollet, T. Miyake, A. Georges, and S. Biermann, *Phys. Rev. B* **80**, 085101 (2009).
- ⁹S. L. Skornyakov, A. V. Efremov, N. A. Skorikov, M. A. Korotin, Yu. A. Izyumov, V. I. Anisimov, A. V. Kozhevnikov, and D. Vollhardt, *Phys. Rev. B* **80**, 092501 (2009).
- ¹⁰S. L. Skornyakov, N. A. Skorikov, A. V. Lukoyanov, A. O. Shorikov, and V. I. Anisimov, *Phys. Rev. B* **81**, 174522 (2010).
- ¹¹S. L. Skornyakov, A. A. Katanin, and V. I. Anisimov, *Phys. Rev. Lett.* **106**, 047007 (2011).
- ¹²D. H. Lu, M. Yi, S.-K. Mo, A. S. Erickson, J. Analytis, J.-H. Chu, D. J. Singh, Z. Hussain, T. H. Geballe, I. R. Fisher, and Z.-X. Shen, *Nature (London)* **455**, 81 (2008); D. H. Lu, M. Yi, S.-K. Mo, J. G. Analytis, J.-H. Chu, A. S. Erickson, D. J. Singh, Z. Hussain, T. H. Geballe, I. R. Fisher, and Z.-X. Shen, *Physica C* **469**, 452 (2009).
- ¹³S. V. Borisenko, V. B. Zabolotnyy, D. V. Evtushinsky, T. K. Kim, I. V. Morozov, A. N. Yaresko, A. A. Kordyuk, G. Behr, A. Vasiliev, R. Follath, and B. Büchner, *Phys. Rev. Lett.* **105**, 067002 (2010).
- ¹⁴C. de la Cruz, Q. Huang, J. W. Lynn, J. Li, W. Ratcliff, J. L. Zarestzky, H. A. Mook, G. F. Chen, J. L. Luo, N. L. Wang, and P. Dai, *Nature (London)* **453**, 899 (2008).
- ¹⁵V. I. Anisimov, A. I. Poteryaev, M. A. Korotin, A. O. Anokhin, and G. Kotliar, *J. Phys.: Condens. Matter* **9**, 7359 (1997); K. Held, I. A. Nekrasov, G. Keller, V. Eyert, N. Blümer, A. K. McMahan, R. T. Scalettar, T. Pruschke, V. I. Anisimov, and D. Vollhardt, *Phys. Status Solidi B* **243**, 2599 (2006).
- ¹⁶W. Metzner and D. Vollhardt, *Phys. Rev. Lett.* **62**, 324 (1989).
- ¹⁷V. I. Anisimov, D. E. Kondakov, A. V. Kozhevnikov, I. A. Nekrasov, Z. V. Pchelkina, J. W. Allen, S.-K. Mo, H.-D. Kim, P. Metcalf, S. Suga, A. Sekiyama, G. Keller, I. Leonov, X. Ren, and D. Vollhardt, *Phys. Rev. B* **71**, 125119 (2005).
- ¹⁸P. Werner, A. Comanac, L. de Medici, M. Troyer, and A. J. Millis, *Phys. Rev. Lett.* **97**, 076405 (2006).
- ¹⁹A. I. Liechtenstein, V. I. Anisimov, and J. Zaanen, *Phys. Rev. B* **52**, 5467R (1995).
- ²⁰P. H. Dederichs, S. Blügel, R. Zeller, and H. Akai, *Phys. Rev. Lett.* **53**, 2512 (1984).
- ²¹O. Gunnarsson, O. K. Andersen, O. Jepsen, and J. Zaanen, *Phys. Rev. B* **39**, 1708 (1989).
- ²²V. I. Anisimov and O. Gunnarsson, *Phys. Rev. B* **43**, 7570 (1991).
- ²³H. J. Vidberg and J. W. Serene, *J. Low Temp. Phys.* **29**, 179 (1977).