

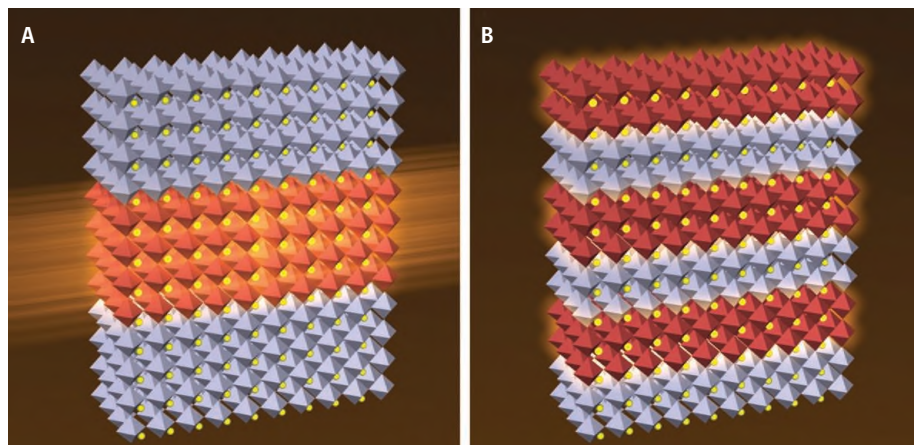
Shedding Light on Oxide Interfaces

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In the stable of solid-state materials, silicon is the workhorse. Its semiconductor properties are actually quite ordinary, but when other materials are harnessed to it to form interfaces, remarkable devices can be made that control the flow of electrical current with low applied voltages. In contrast, some transition metal oxides would be the stable's racehorses. Remarkable functional behaviors of this group include the ability to change from a metal to an insulator with a slight change in temperature, unusual magnetic properties, and even high-temperature superconductivity (1). Further novel properties are expected to emerge at interfaces created between transition metal oxides that already exhibit functional behavior, and these properties could be tuned through small changes in composition or by simply applying a bias voltage (2). However, like temperamental racehorses, the interfaces in these complex oxides can be more difficult to control than those formed by silicon. On page 937 of this issue, Boris *et al.* (3) report prog-

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The properties of a metallic oxide can be altered when it is confined as an ultrathin layer by layers of an insulating oxide.



Thinner is different. Artistic illustration of two different electronic phases in heterostructures of LaNiO₃ (red), which is metallic as a bulk material, confined by LaAlO₃ (gray), which is a bulk insulator. **(A)** A confined LaNiO₃ layer that is four unit cells in thickness still shows metallic behavior. **(B)** When the thickness is reduced to two unit cells, a temperature-driven metal-insulator transition emerges.

ress toward this goal by confining an ultrathin layer of lanthanum nickelate (LaNiO₃), normally a paramagnetic metal, between insulating lanthanum aluminate (LaAlO₃) layers, which leads to changes in its properties.

The difference between semiconductors such as silicon and exotic transition metal oxides lies in the nature of the interactions between the constituent electrons. The electronic states of semiconductors are well

described by conventional, single-particle band theory, in which individual electrons act independently. Contact between energy bands at interfaces can bend the bands or create energy barriers, which in turn can generate quasi-electric fields (which act on the charge carriers differently from true electric fields) or cause accumulation of charge carriers. In the complex oxides, however, there are strong correlations between the tightly

bound transition metal d electrons, and single electrons are no longer independent but influence the behavior of all other electrons in the crystal. The large electrical polarizability of oxygen also enhances the response of these materials to electric fields causing ferroelectric (a net electric dipole moment) or multiferroic (both ferroelectric and ferromagnetic) properties (4).

Boris *et al.* used pulsed laser deposition to grow superlattices—precise numbers of oxide layers with atomically sharp interfaces (see the figure). Samples with four-unit-cell thick LaNiO_3 layers grown between LaAlO_3 layers were metallic at temperatures from 8 to 300 K, just like bulk LaNiO_3 . Samples with only two layers of LaNiO_3 , in which each LaNiO_3 layer is next to a LaAlO_3 layer, showed a metal-to-insulator transition with subsequent magnetic ordering as the temperature was lowered.

Such effects might be mainly the result of strain in the layers, which can be induced by the substrate used for growth. To check whether this was the case, Boris *et al.* grew the same layer structures on strontium titanate (SrTiO_3), which has a larger lattice constant than the superlattice, and on lanthanum strontium aluminate (LaSrAlO_4), in which the lattice constant is smaller. The same behavior was seen in both cases, except for a shift of transition temperature (100 and 150 K, respectively), ruling out strain as the cause.

To demonstrate further that these effects were intrinsic to the superlattice structure, Boris *et al.* applied state-of-the-art surface probes. They used ellipsometry to measure the changes in the sample's electrical conductivity. Ellipsometry is not influenced by extrinsic impurities, particularly interdiffusion of ions between the layers, nor is it influenced by misfit dislocations resulting from inexact matching of the lattice constants at the interface, which can occur in oxide superlattices. They also used low-energy muon spin rotation to identify a change in the magnetic order accompanying the transition from the metallic to the insulating state. Here, spin-polarized muons with well-defined energies are implanted within the oxide heterostructure, where they align their magnetic moments parallel to those of the surrounding electrons. The measured directions of the subsequent positron decay products yield the local magnetic order of the electronic phase. The bilayered LaNiO_3 samples showed a clear magnetic transition, but the superlattices with thicker layers did not.

Deposition techniques such as pulsed laser deposition or molecular beam epitaxy

are now well developed, so nearly defect-free superlattices with atomically sharp interfaces can be routinely grown. Unanticipated phenomena such as the metal-insulator transition described here, or conductivity at interfaces between the insulating oxides such as LaAlO_3 and SrTiO_3 (5) or interfacial superconductivity (6), are now regularly reported. However, there is still much work to be done before oxide materials reach the level of sophistication achieved in semiconductor heterostructures. Improved control and understanding of the role of defects is necessary, as well as systematic incorporation of multiple interfaces with different electronic properties (2).

From a practical point of view, demonstrations of integration with conventional semiconductors would be helpful, as would a detailed understanding of the behavior at interfaces with metallic electrodes (7). On the theoretical front, the aspects that make complex oxides desirable—in particular the strong correlations, the large polarizability, and the

sensitive dependence on crystal chemistry and structure—also make them challenging to describe accurately. Improved techniques must combine many-body physics methods for describing strong correlations with computational materials methods, such as density functional theory, that can account for chemistry and structure. With these developments, which the community is poised to make over the next few years, true predictive capability and layer-by-layer construction of designer oxide superlattices should be achievable.

References

1. D. G. Schlom, L.-Q. Chen, X. Pan, A. Schmehl, M. A. Zurbuchen, *J. Am. Ceram. Soc.* **91**, 2429 (2008).
2. J. Mannhart, D. G. Schlom, *Science* **327**, 1607 (2010).
3. A. V. Boris *et al.*, *Science* **332**, 937 (2011).
4. N. A. Spaldin, S. W. Cheong, R. Ramesh, *Phys. Today* **63**, 38 (2010).
5. A. Ohtomo, H. Y. Hwang, *Nature* **427**, 423 (2004).
6. A. Gozar *et al.*, *Nature* **455**, 782 (2008).
7. M. Stengel, N. A. Spaldin, *Nature* **443**, 679 (2006).

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